

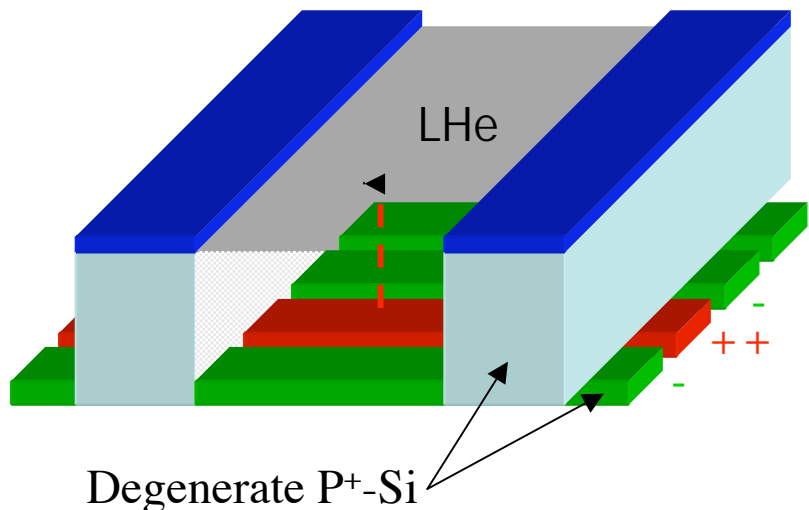
Electron Spin Qubits on Liquid Helium

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Princeton University

1. Electron spins as qubits
2. Quantum dot plans for single spin measurements
3. Clocking electrons on the surface of helium
4. Photoemission electron source

Helium channels for electron transport

- For storing and moving qubits, we need a fairly high qubit density – electrons a few μm apart
- To electrostatically control individual electrons, the electrodes must be no farther than a few μm from them
- The channels are a convenient way to produce helium films a few μm deep, in a controlled manner



- The electron motion is controlled with electrodes below the channels
- This is the structure of a buried-channel Si CCD
- Channel width and gate period = 4 μm ; depth = 2 μm

Spin Decoherence Times

- Can move qubits, so normally keep them far apart (dipole-dipole at $10\mu\text{m} \sim 10^4 \text{ sec}$)
- Residual spin-orbit interaction ($\mathbf{v} \times \mathbf{E}$)
 - $v \sim 10^6 \text{ cm/sec}$, $E \sim 10^4 \text{ V/cm} \Rightarrow B_{\text{effective}} = B_{\parallel} \sim 10^{-7} \text{ T}$
 - B_{\parallel} changes direction as electron scatters, so

$$1/T_2 = \frac{1}{2} \gamma^2 \overline{B_{\parallel}^2} \frac{\tau}{1 + \omega_0^2 \tau^2}$$
 - For $\tau \sim 100 \text{ nsec} \Rightarrow T_2 \sim 10^6 \text{ sec}$
- ^3He impurities on ^4He surface ($\sim 0.01/\mu\text{m}^2$) with motional narrowing $\Rightarrow T_2 \sim 10^5 \text{ sec}$
- Johnson noise currents in gate layers:

$$T_2 = \frac{64\pi}{\gamma^2 \mu_0^2 \sigma} \cdot \frac{d(d+t)}{t} \cdot \frac{1}{kT + \frac{3}{4} \hbar \omega_0 \coth(\frac{\hbar \omega_0}{2kT})}$$
 - d = distance to metal, t = metal film thickness
 - For degenerately-doped Si, $d \sim 2\mu\text{m}$, $T \sim 30 \text{ mK} \Rightarrow T_2 \sim 10^5 \text{ sec}$
 - On thin helium (quantum dot & 2-qubit gate), degenerately-doped Si electrode, $d \sim 400\text{\AA}$, $t \sim 1000\text{\AA} \Rightarrow T_2 \sim 1000 \text{ sec}$ at 30mK
- Fluctuations of local spin density in nearby metal
 - Thin helium, degenerately-doped Si electrode (p-type, T_1 for holes $\sim 30\text{ps}$), $d \sim 400\text{\AA}$, $t \sim 1000\text{\AA} \Rightarrow T_2 \sim 2000 \text{ sec}$ at 30mK
- Nuclear spins in gates/channels (use degenerately boron-doped ^{28}Si)
 - Free carriers rapidly relax nuclear spins $\Rightarrow > 10^4 \text{ sec}$
- Paramagnetic defects?
 - Few defects with unpaired spins in Si, and those relaxed rapidly by free carriers

Quantum dots to measure spins

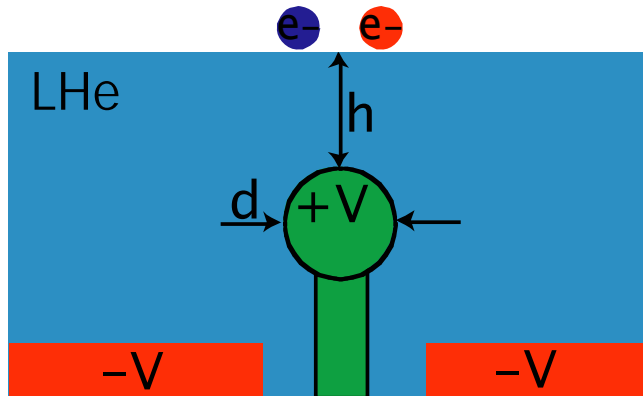
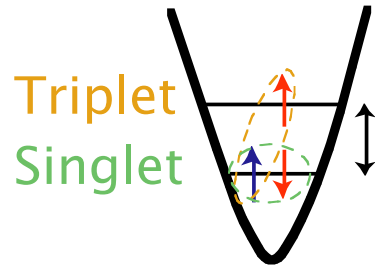
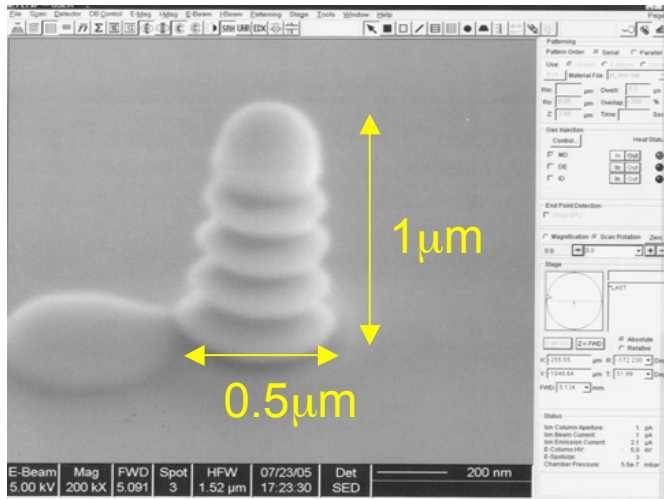
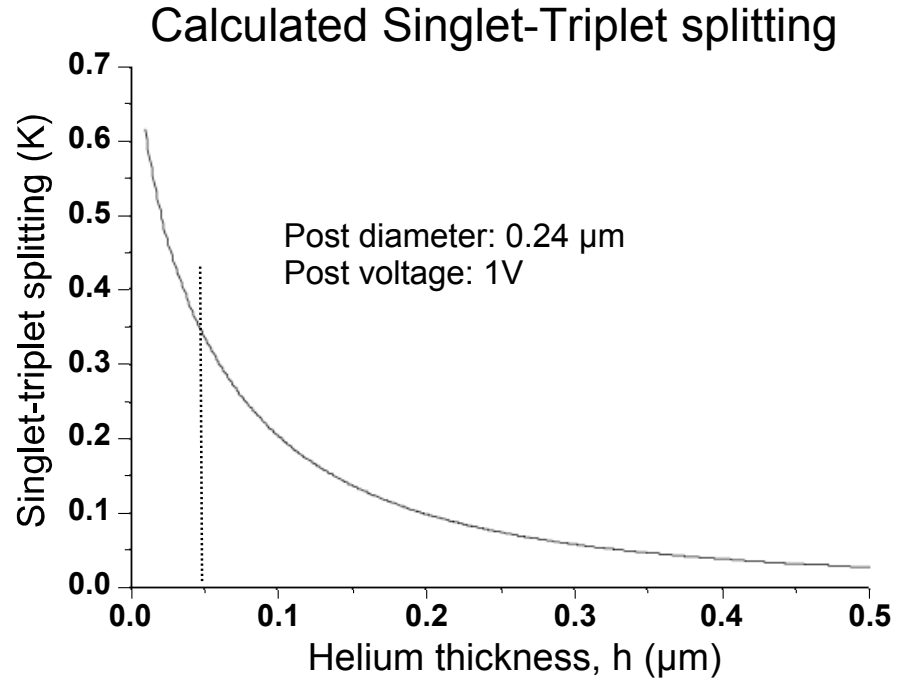


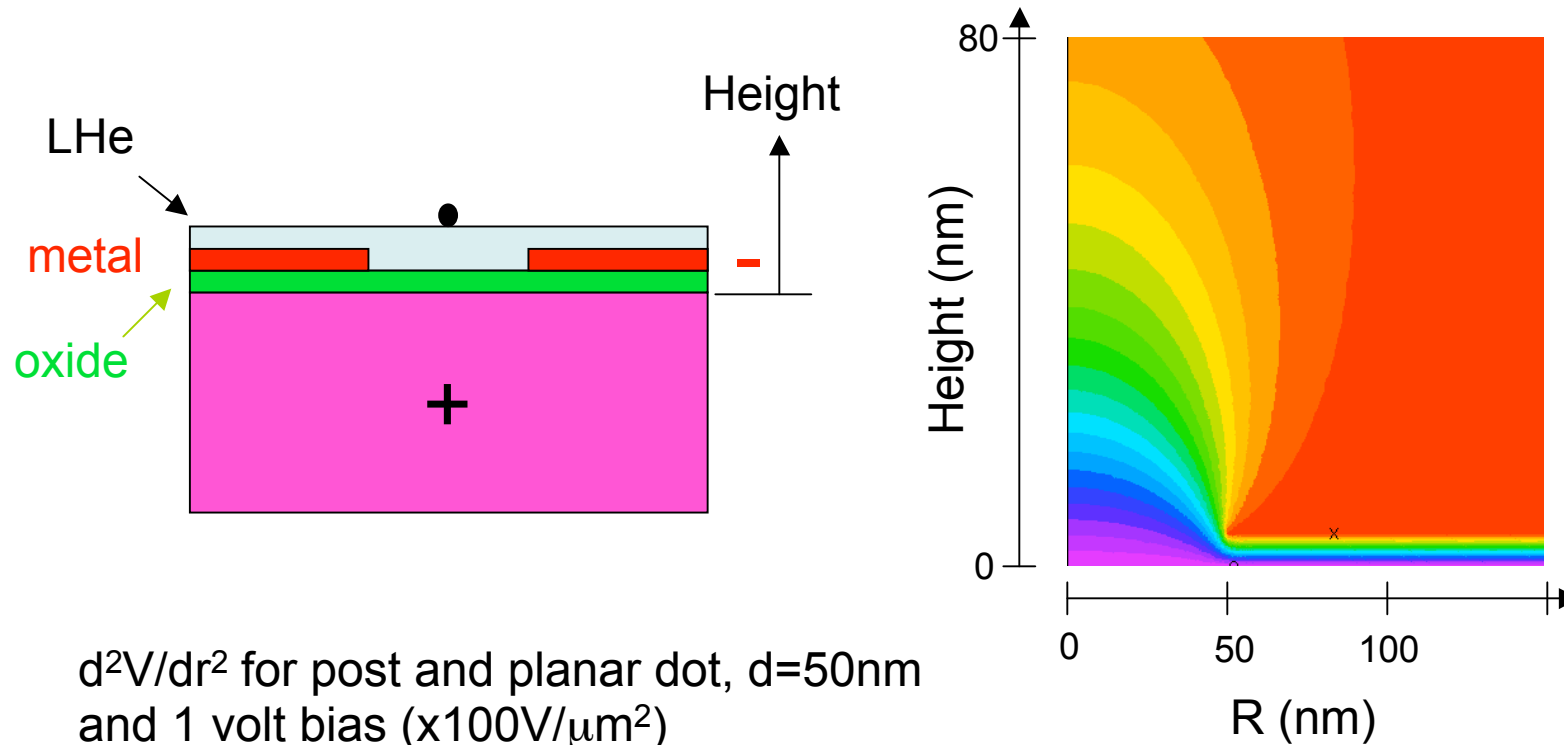
Diagram of a post and its harmonic potential



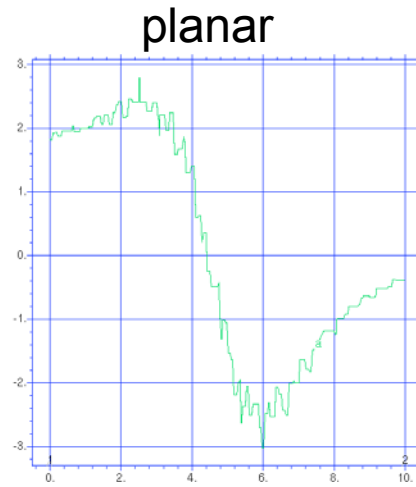
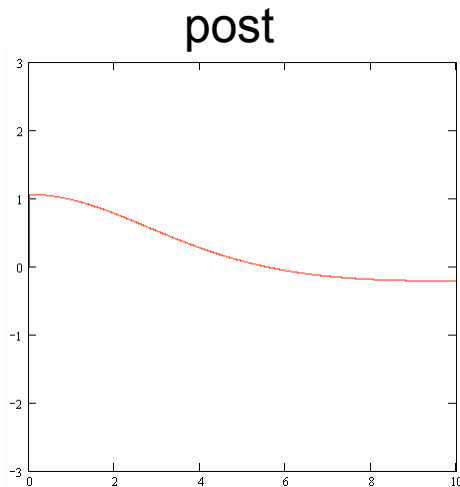
Platinum post made with FIB

A post creates a harmonic potential that splits the singlet and triplet energy levels.

Planar quantum dot



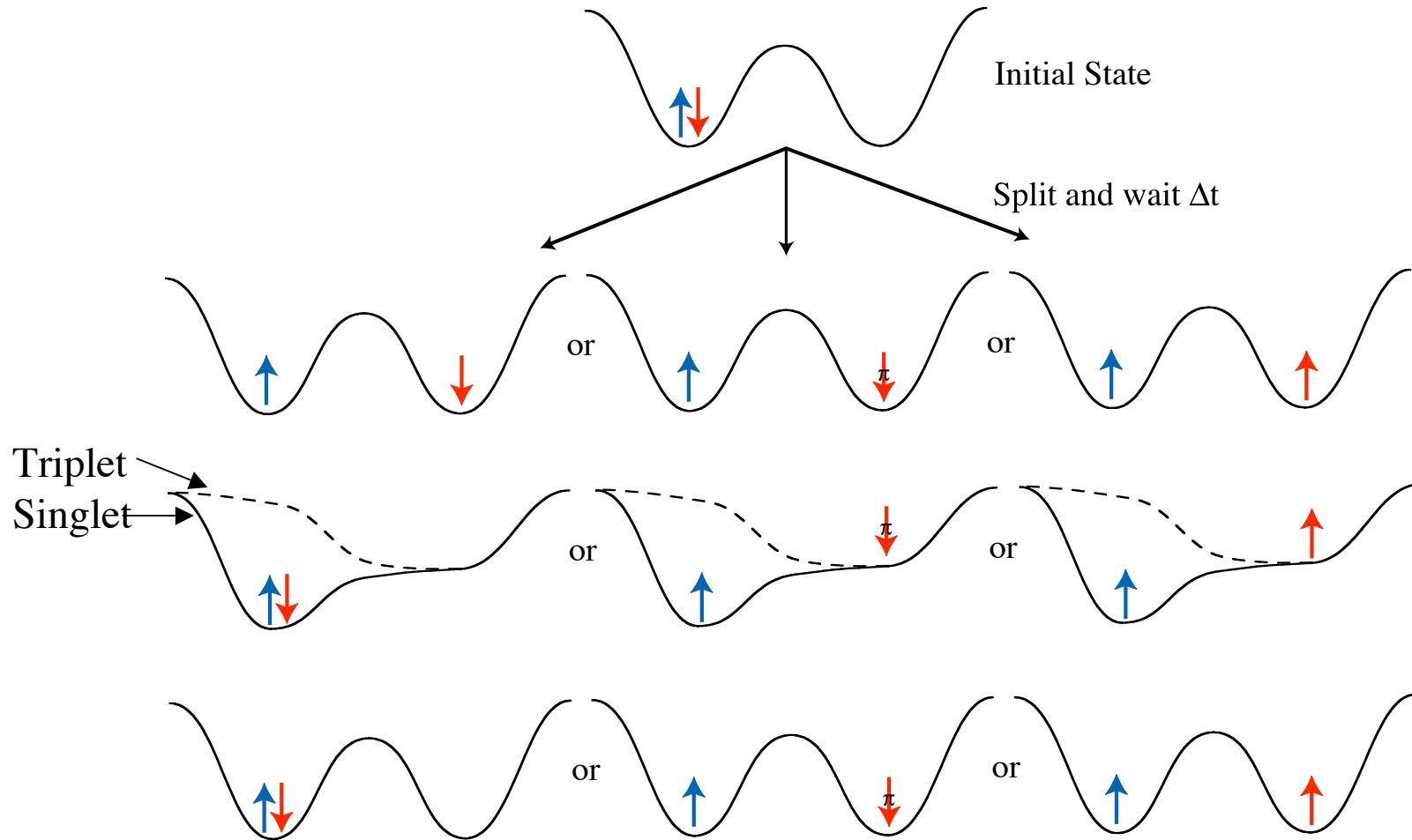
d^2V/dr^2 for post and planar dot, $d=50\text{nm}$ and 1 volt bias ($\times 100\text{V}/\mu\text{m}^2$)



Planar dot looks to be easier to make than the post, and should give $\sim 2\times$ larger level spacing \Rightarrow we'll try this first

Spin Coherence - Measurement Plans

(like in GaAs quantum dots)



Planar double quantum dot (Dipole-Dipole CNOT)

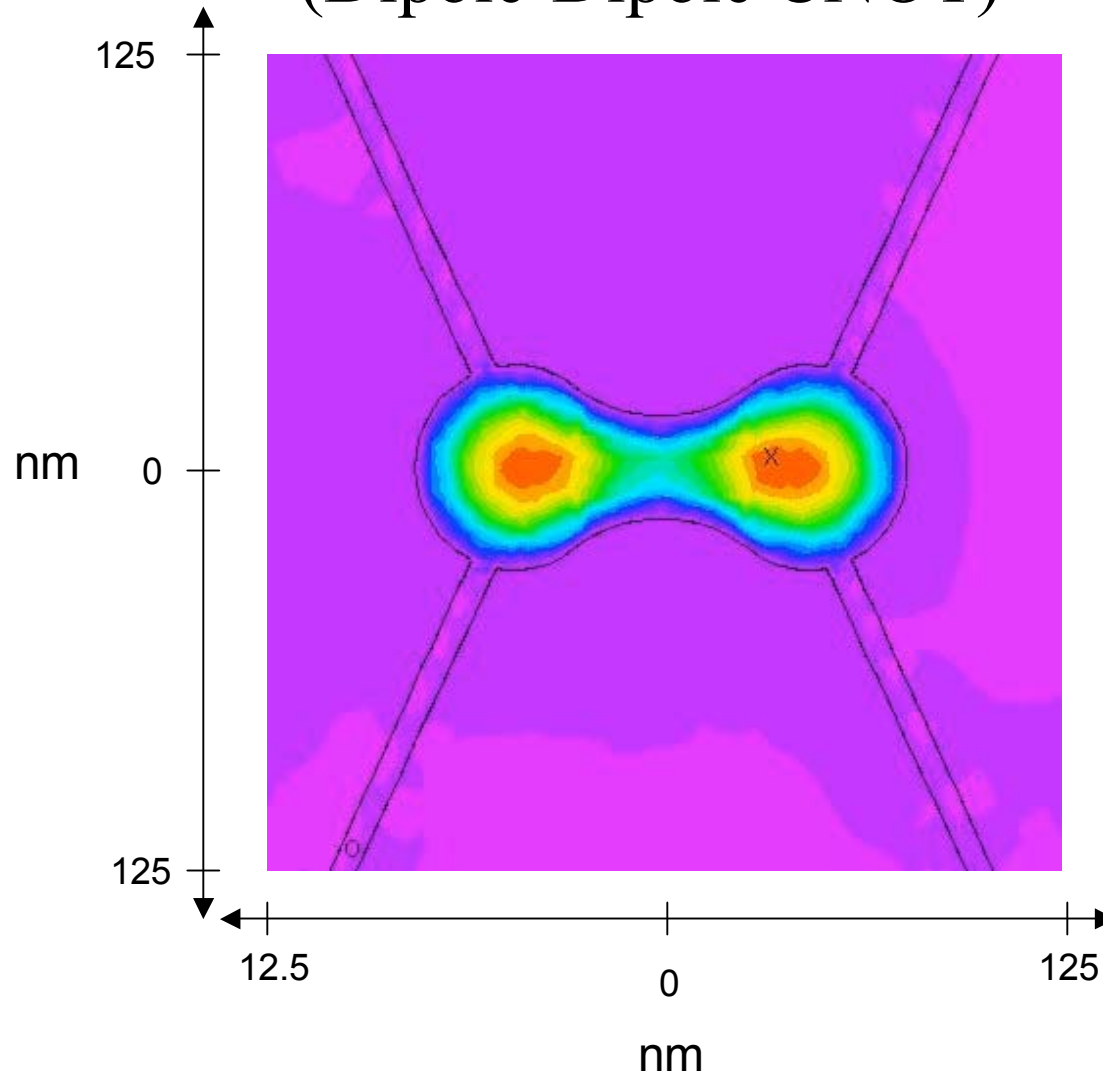
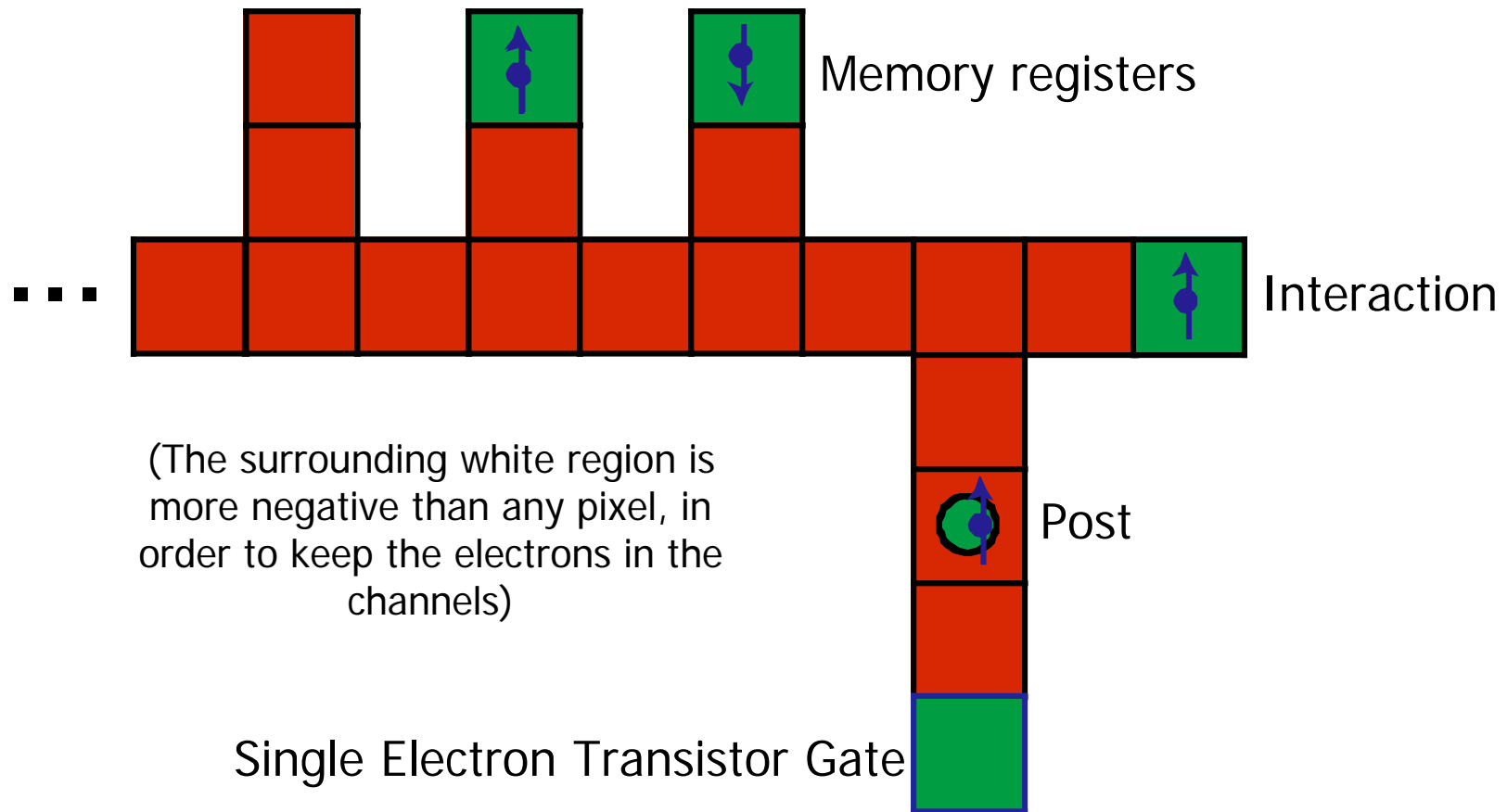
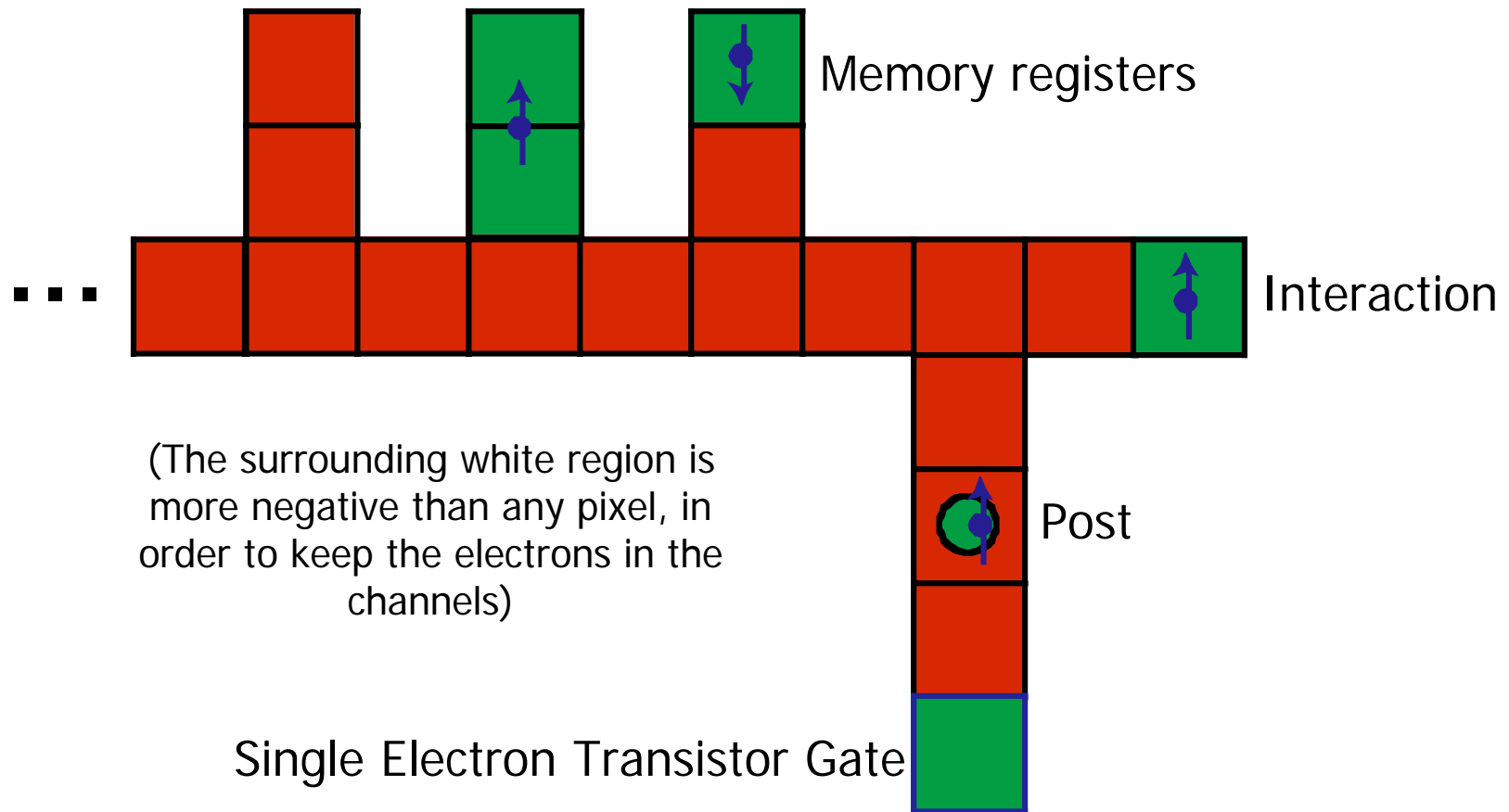


Illustration of Pixel Sequence



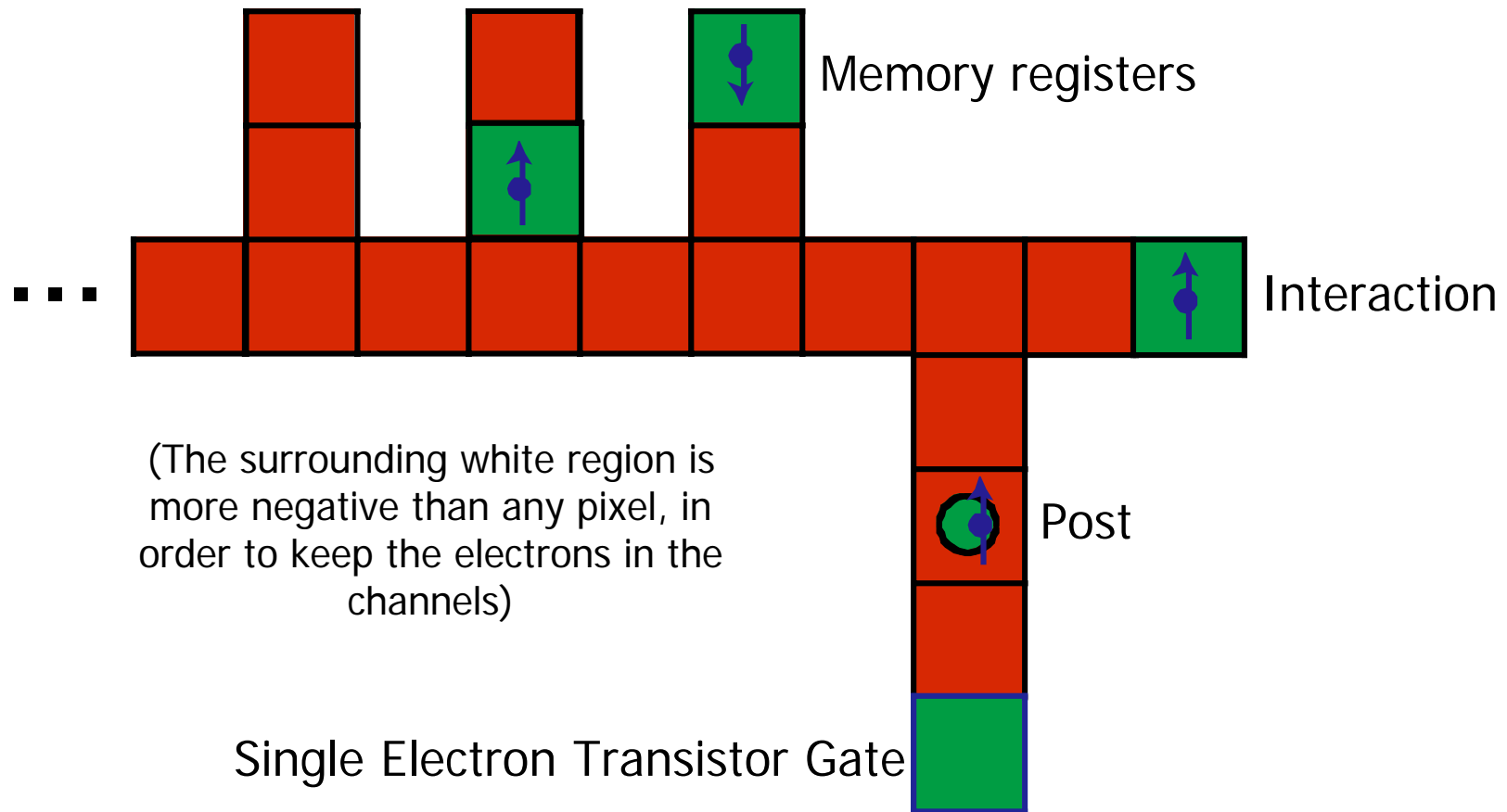
- An electron is taken out of memory, brought to interact with another electron and is then read.

Illustration of Pixel Sequence



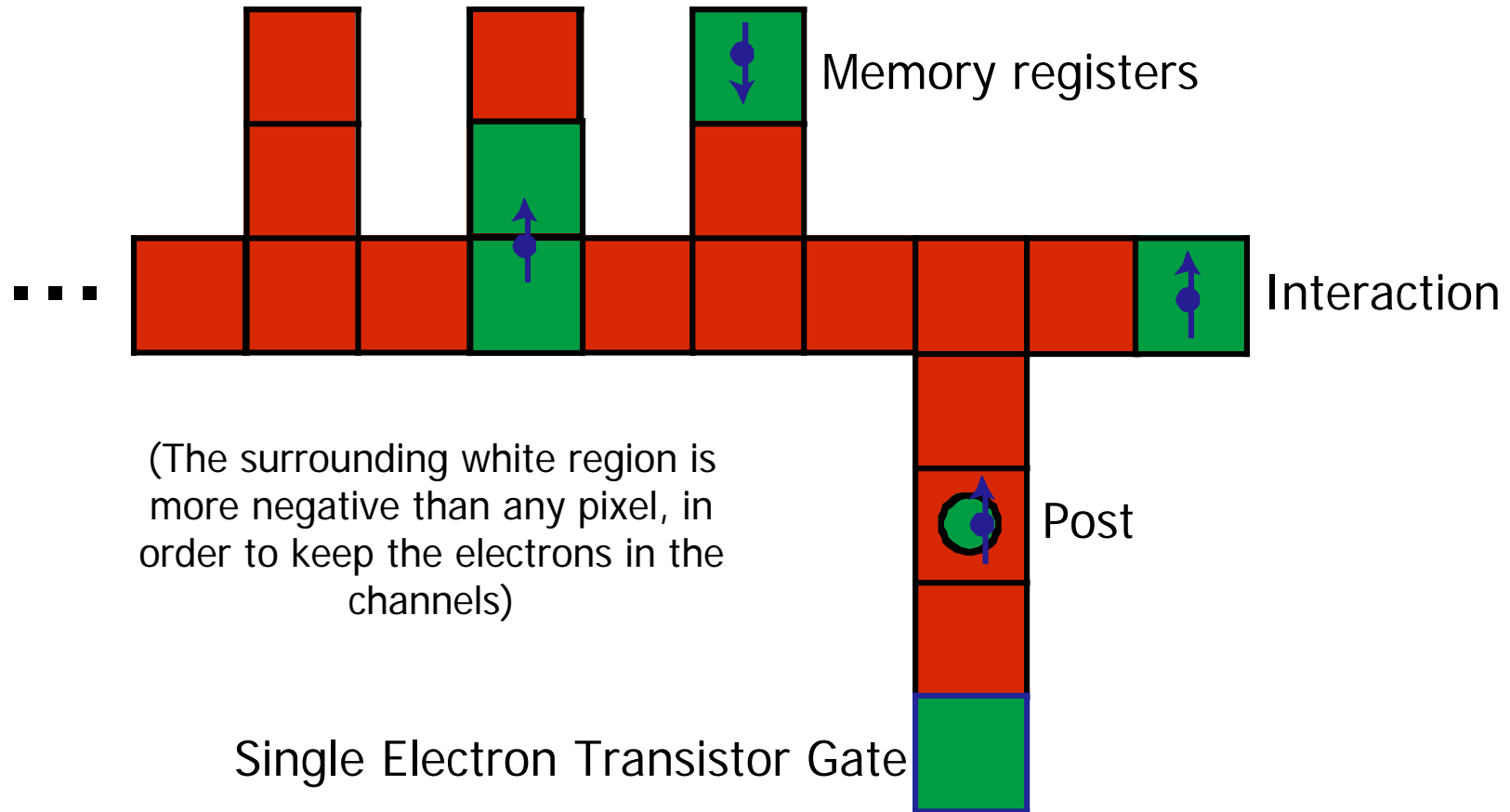
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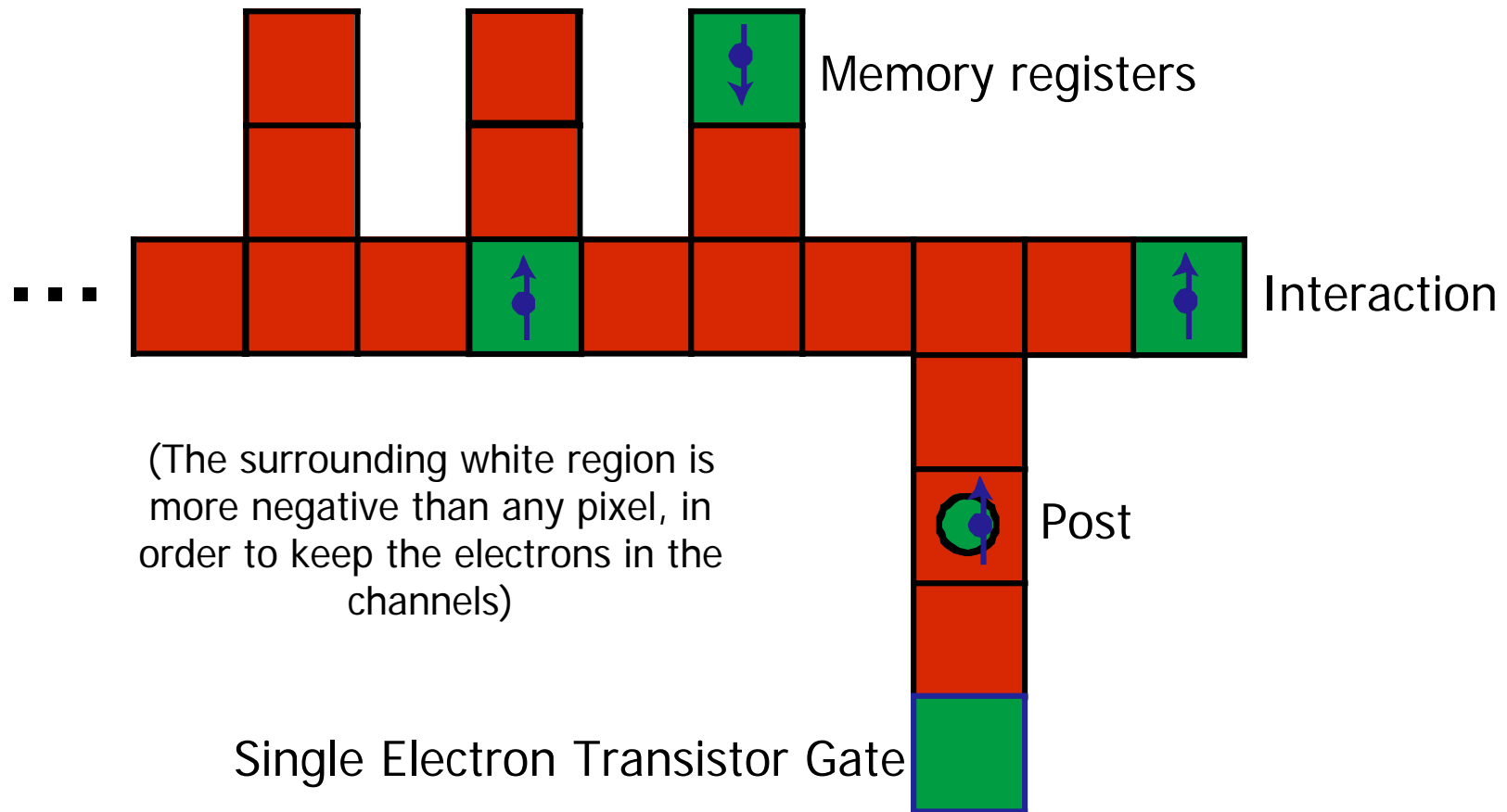
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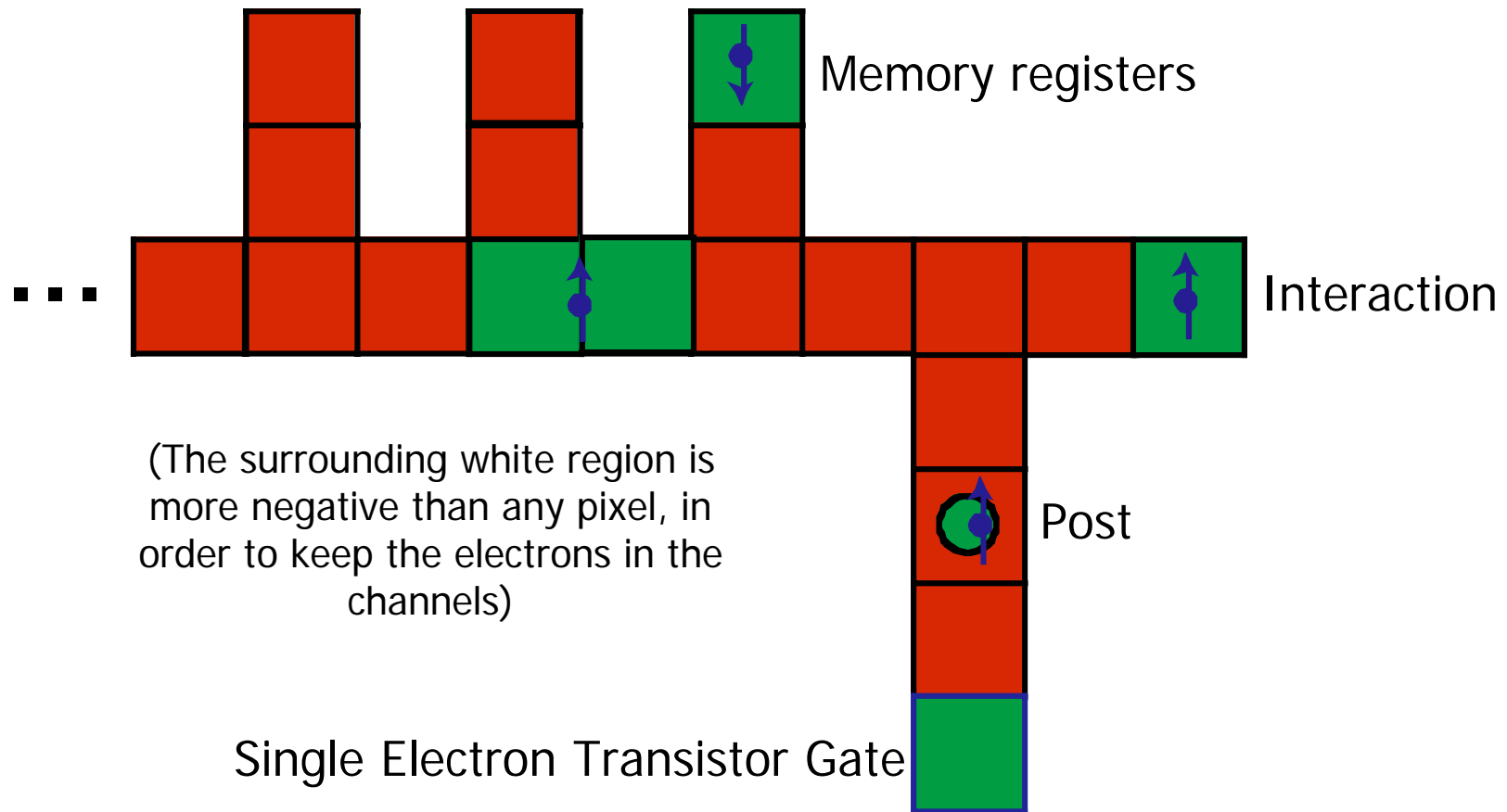
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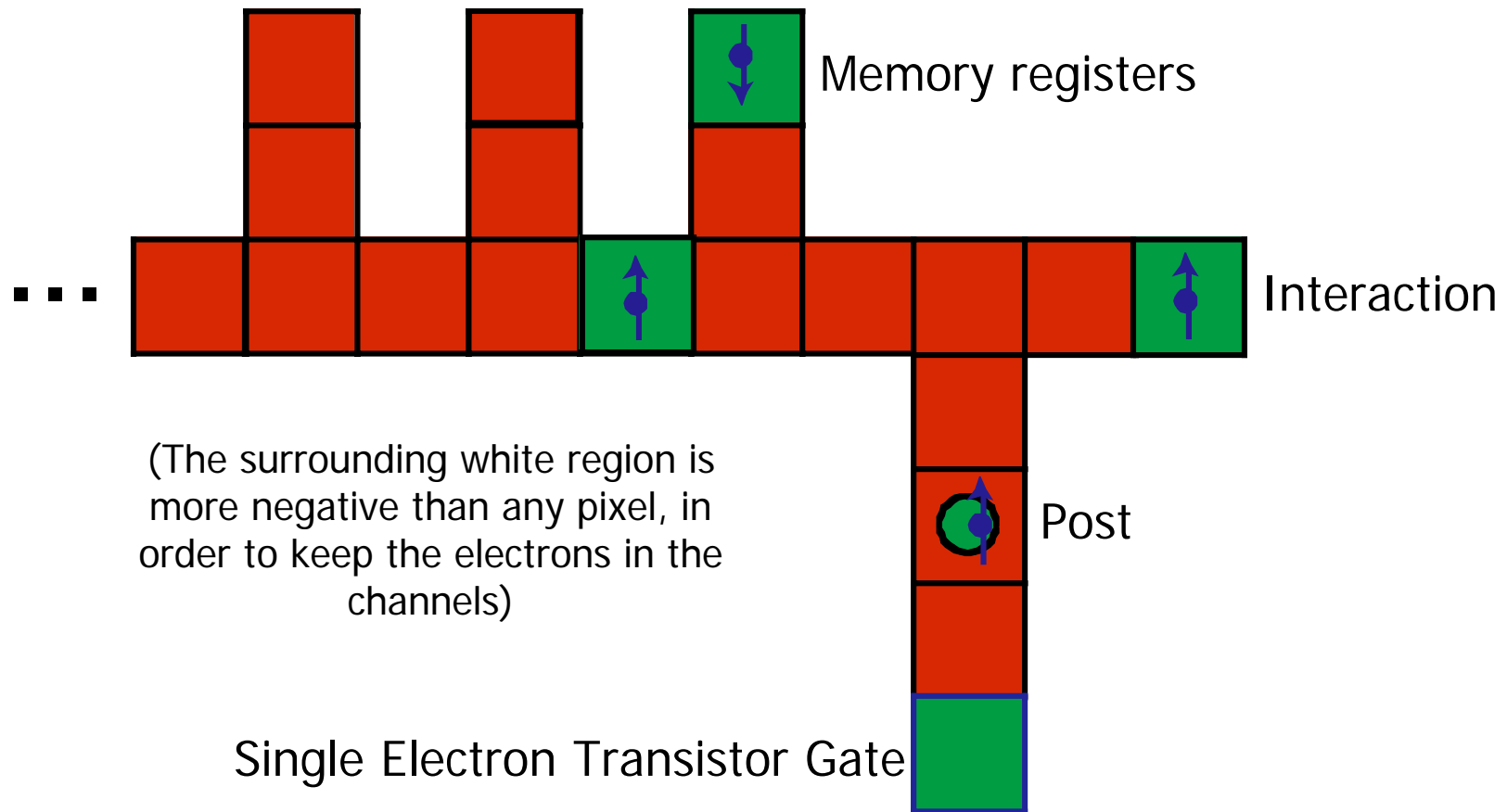
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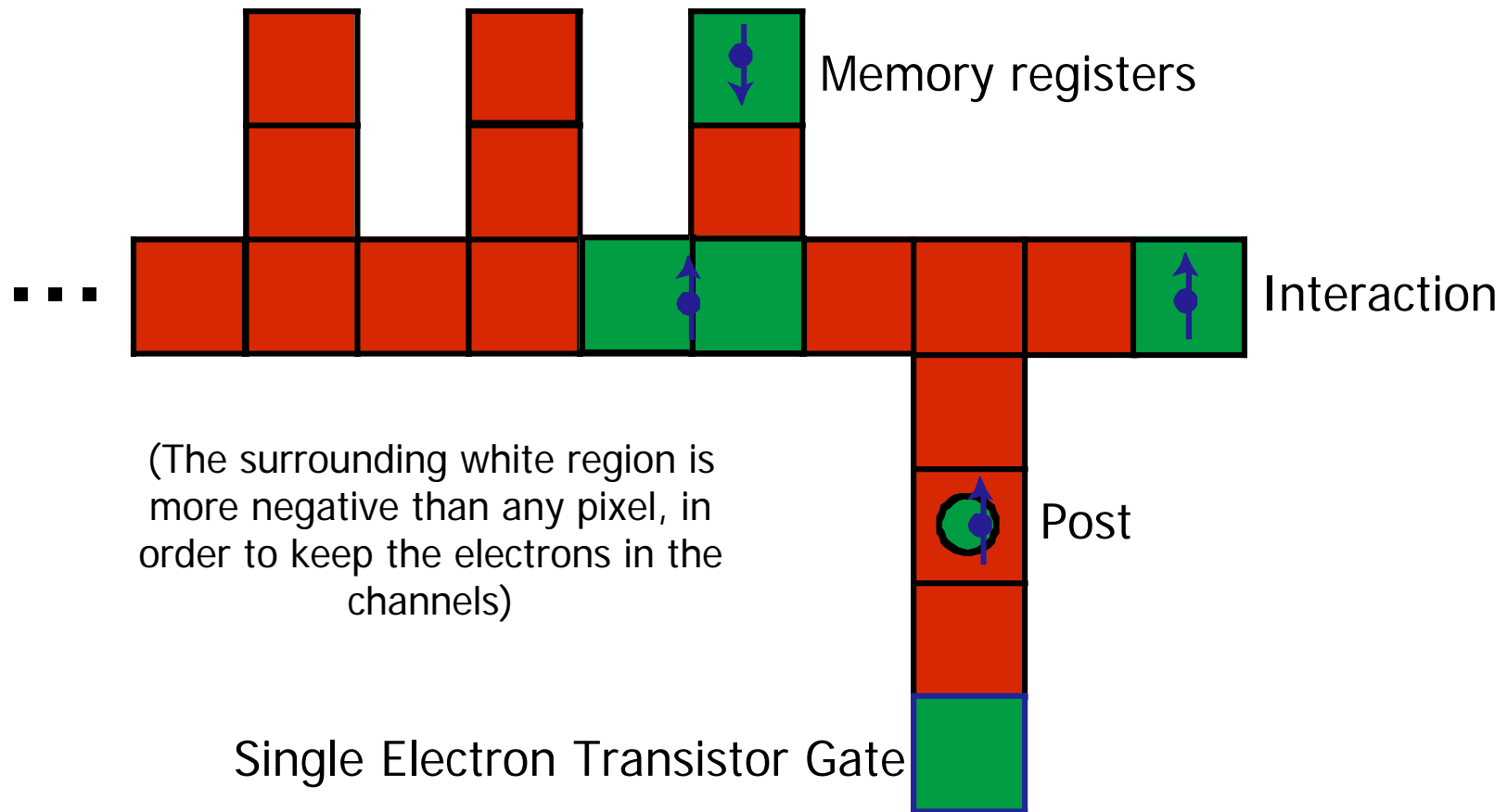
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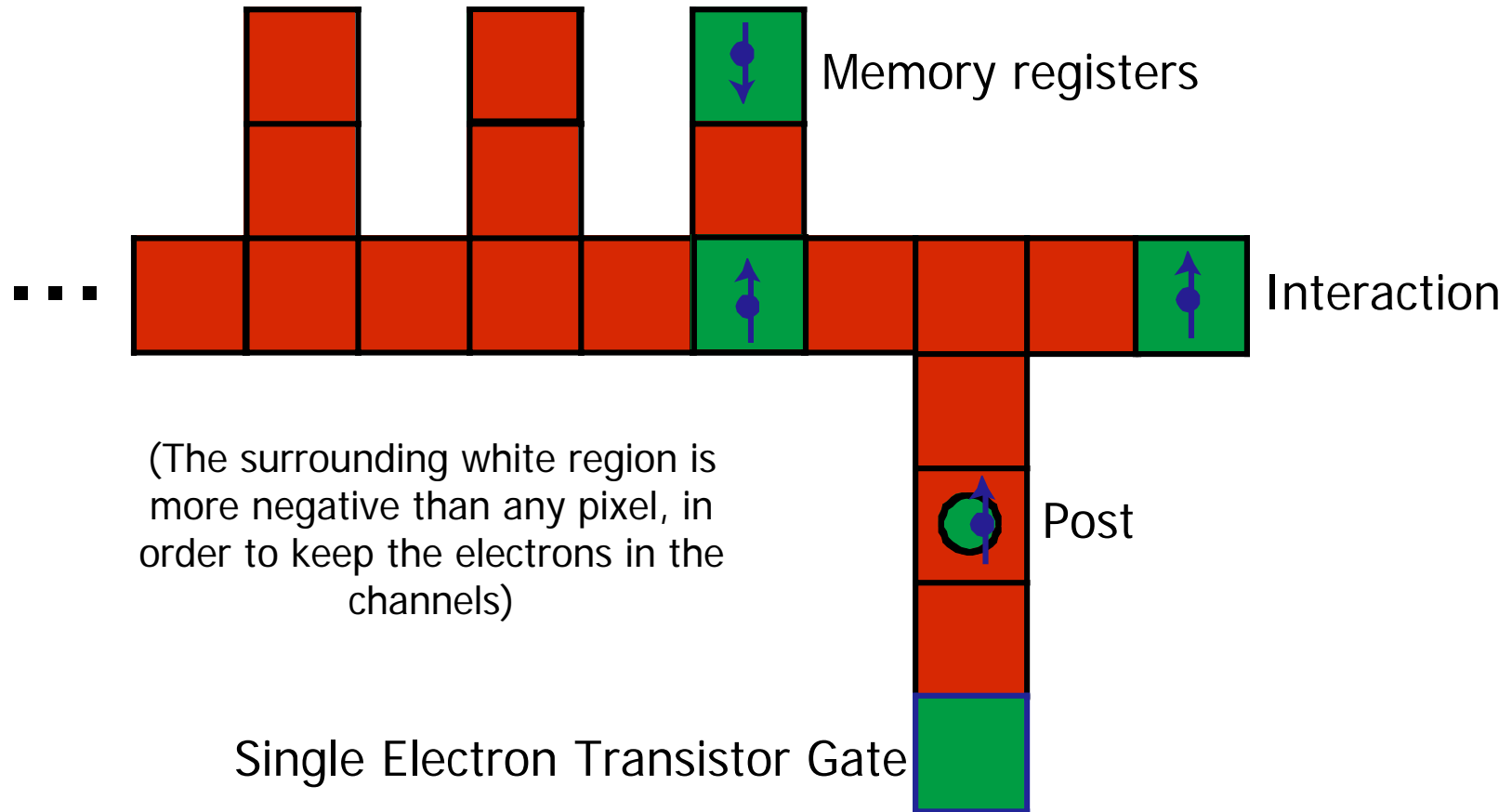
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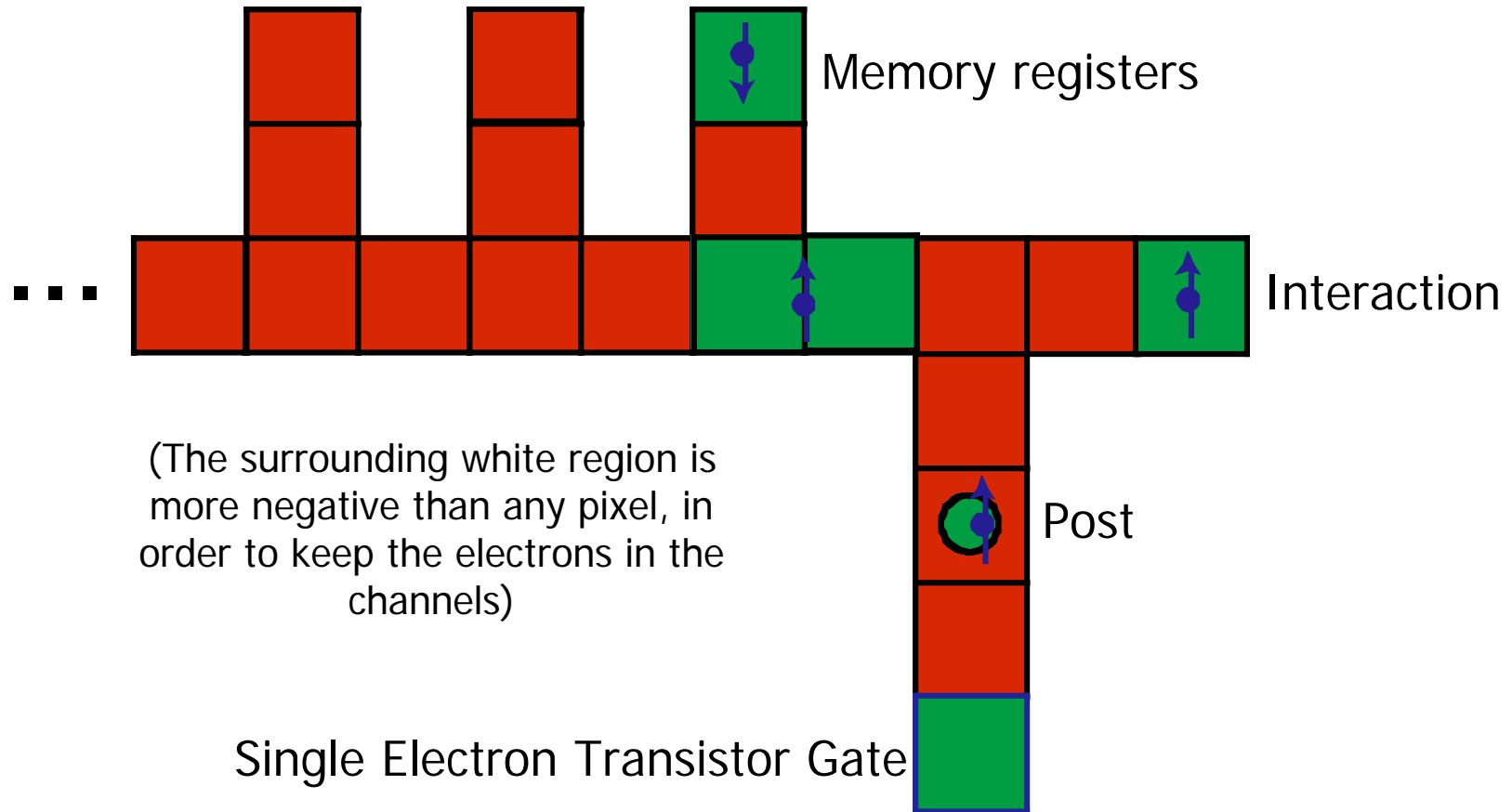
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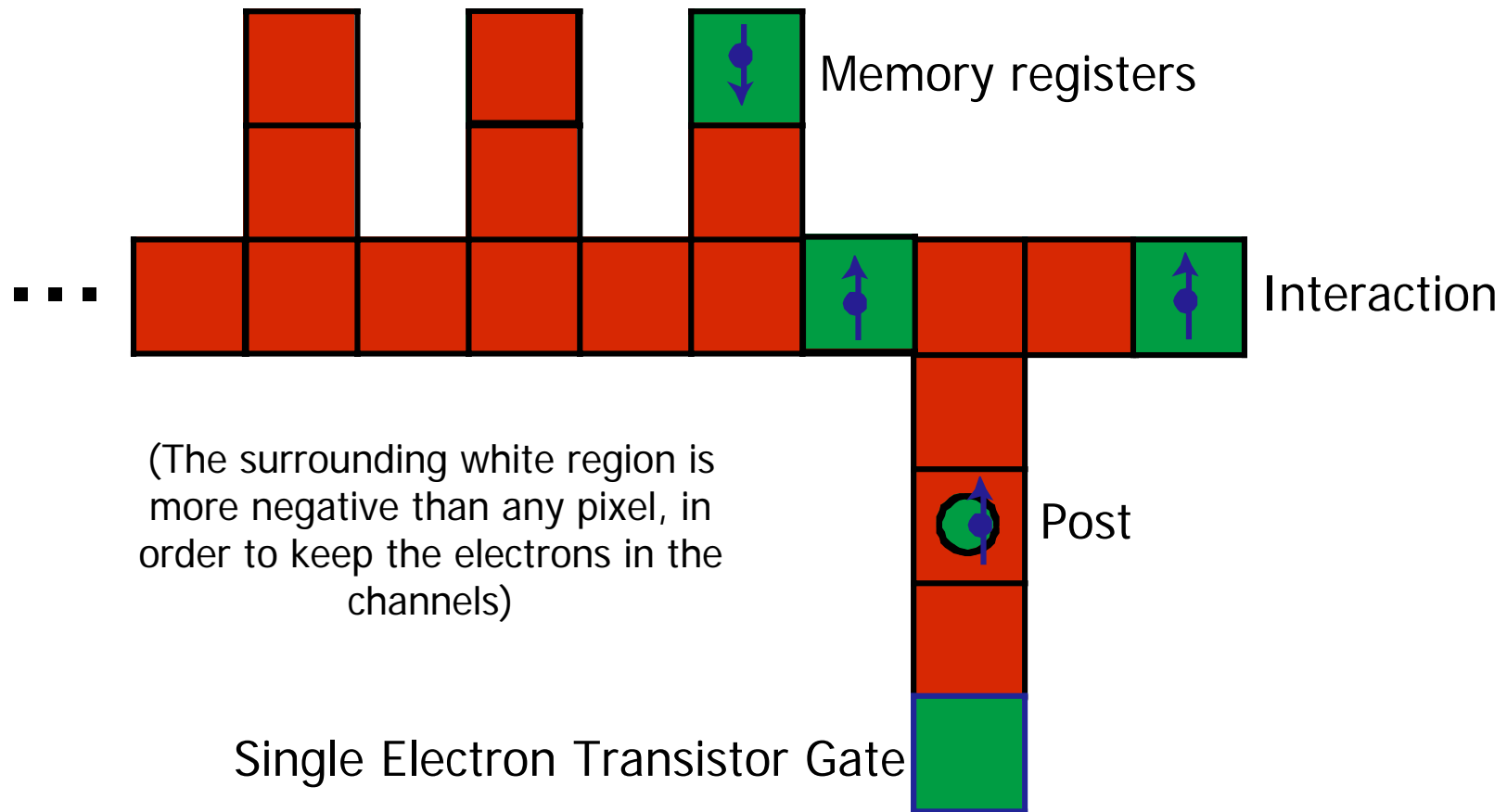
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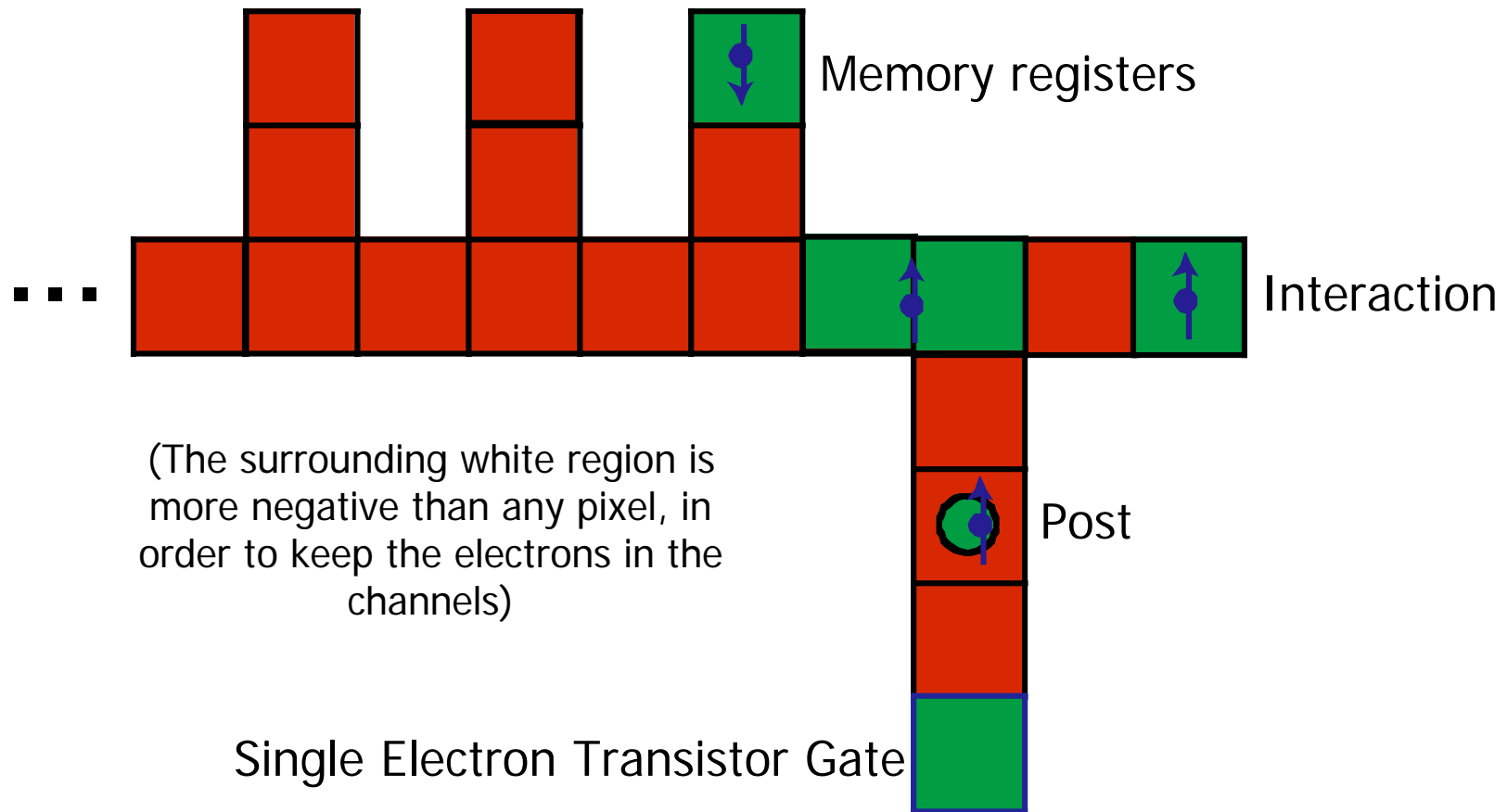
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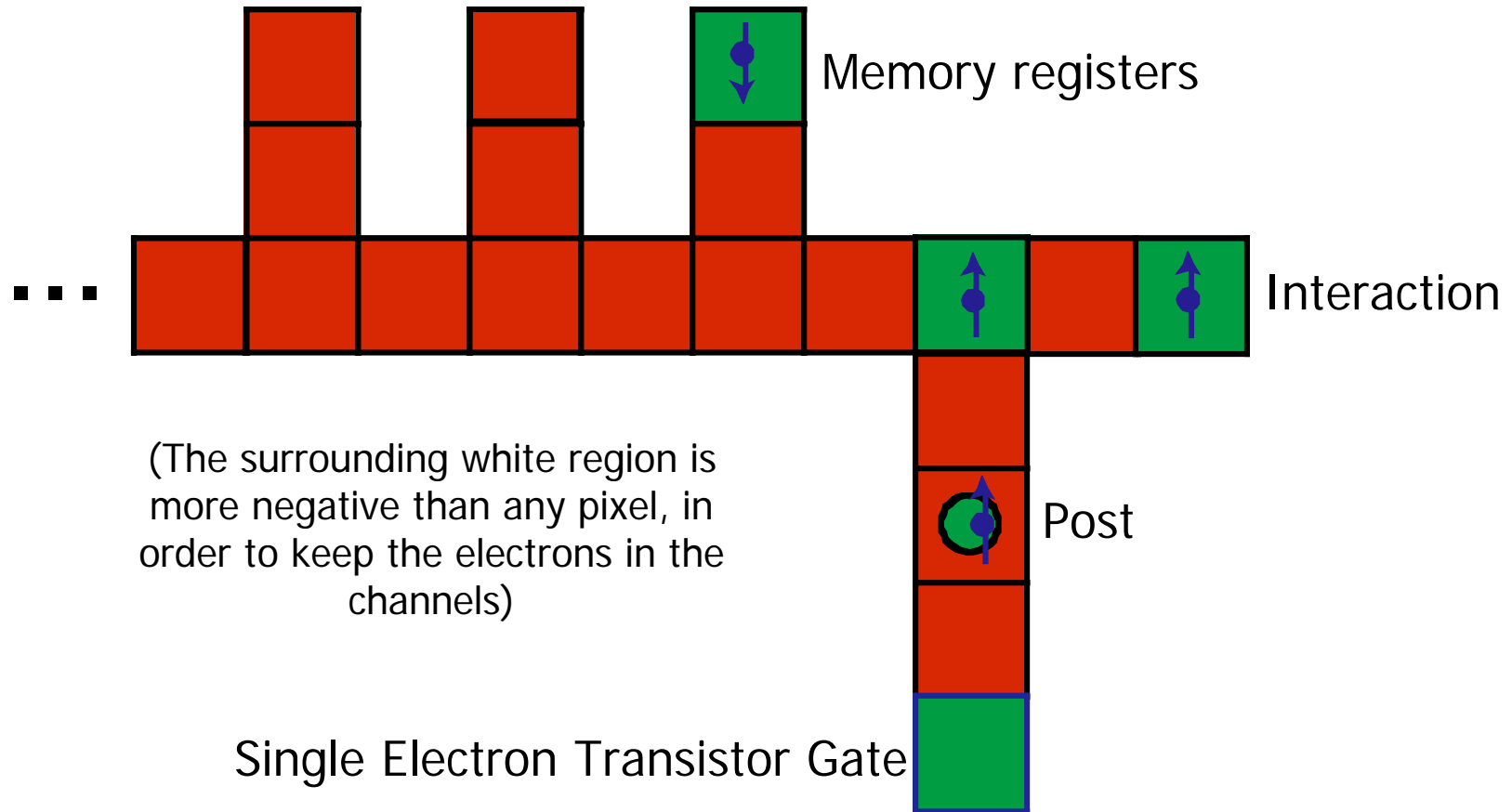
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Illustration of Pixel Sequence



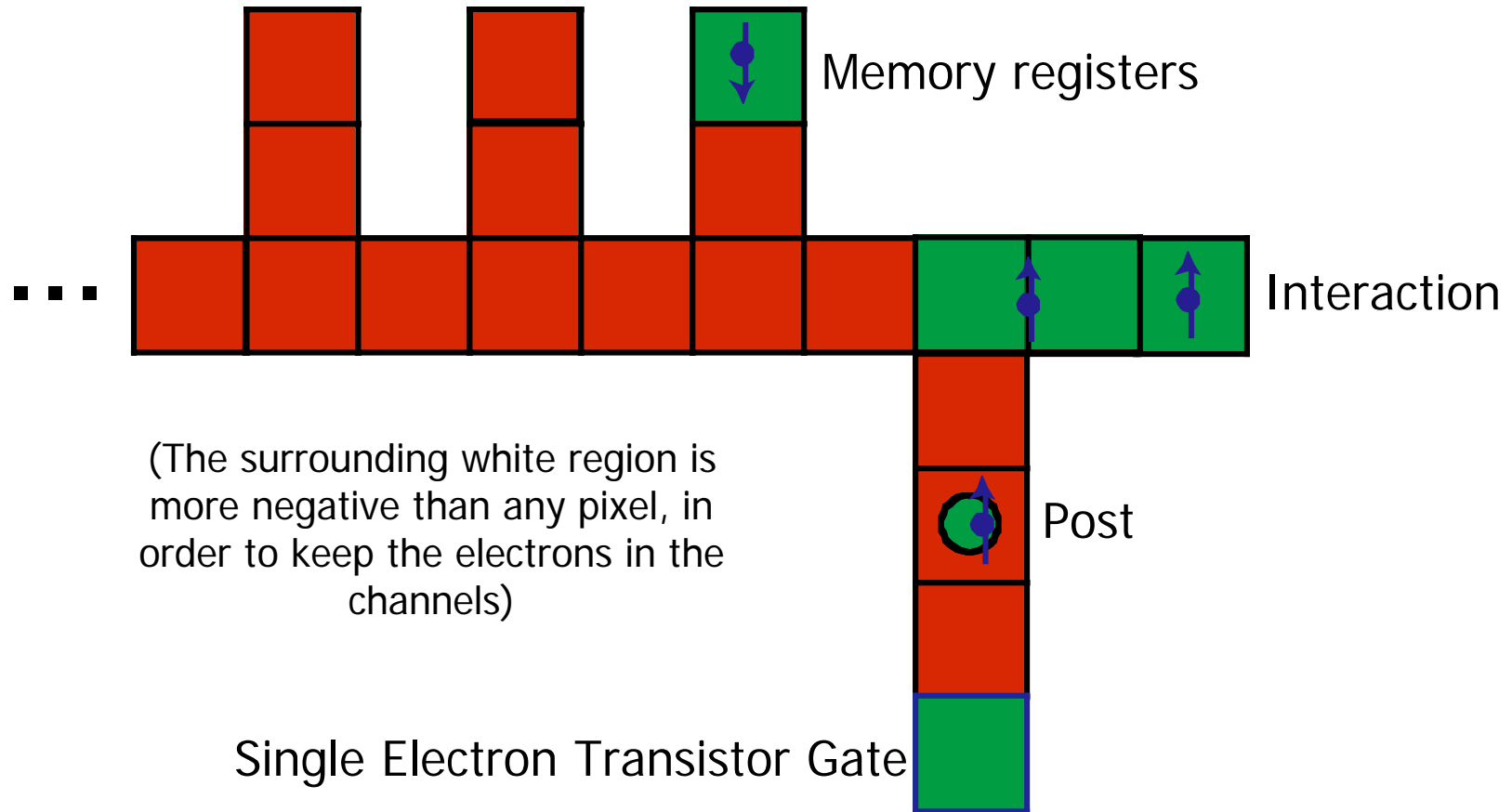
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Illustration of Pixel Sequence



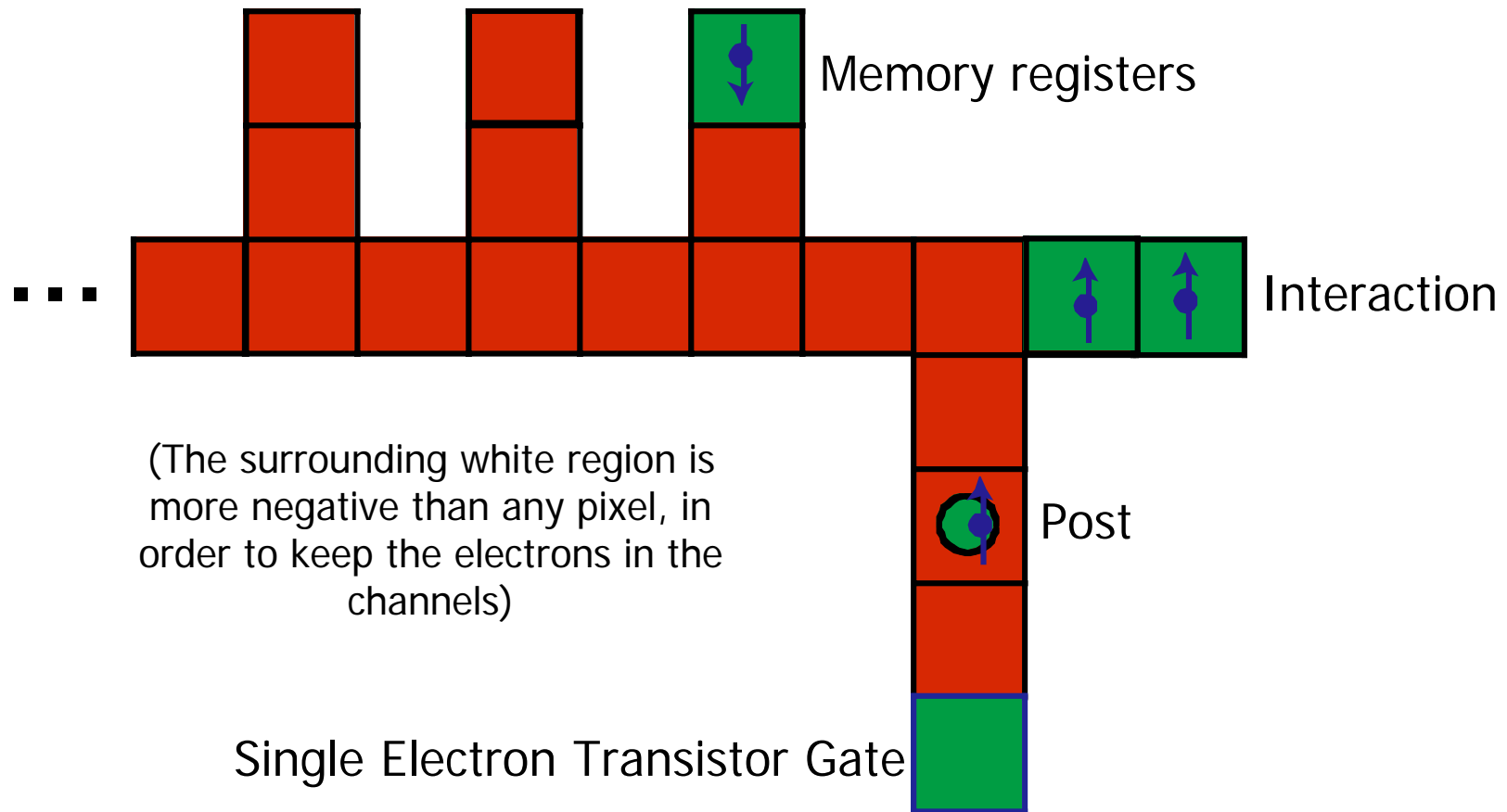
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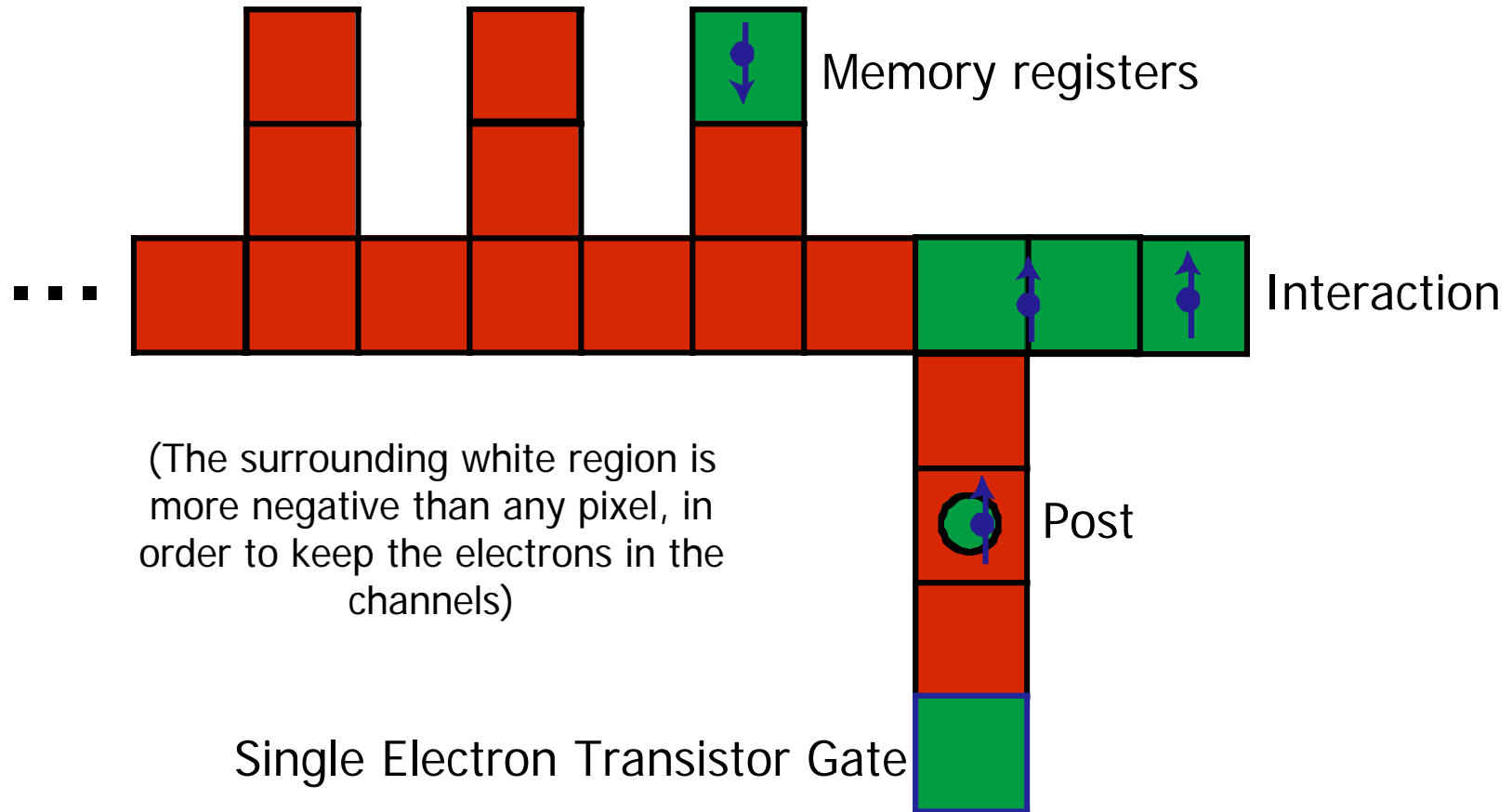
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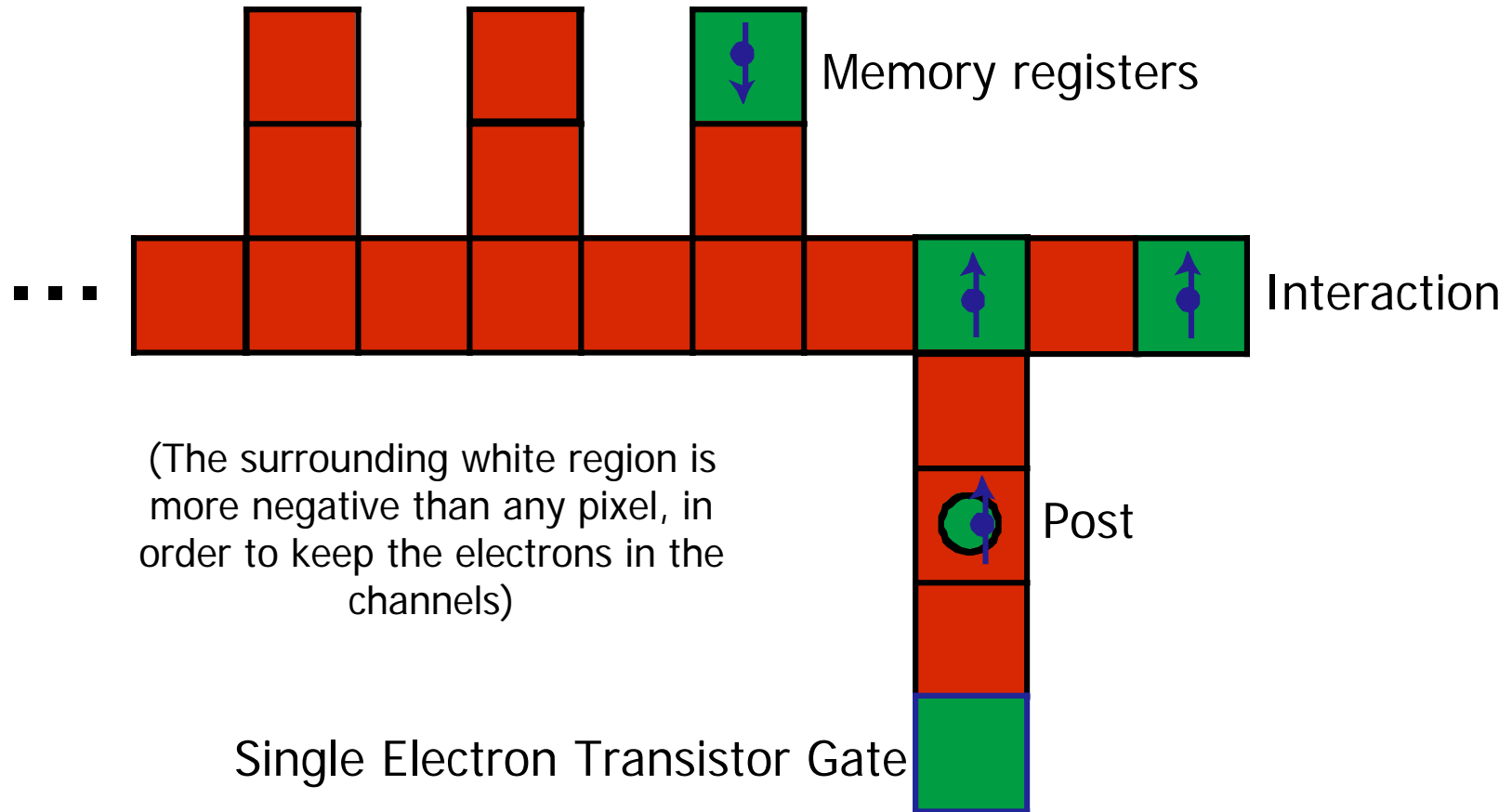
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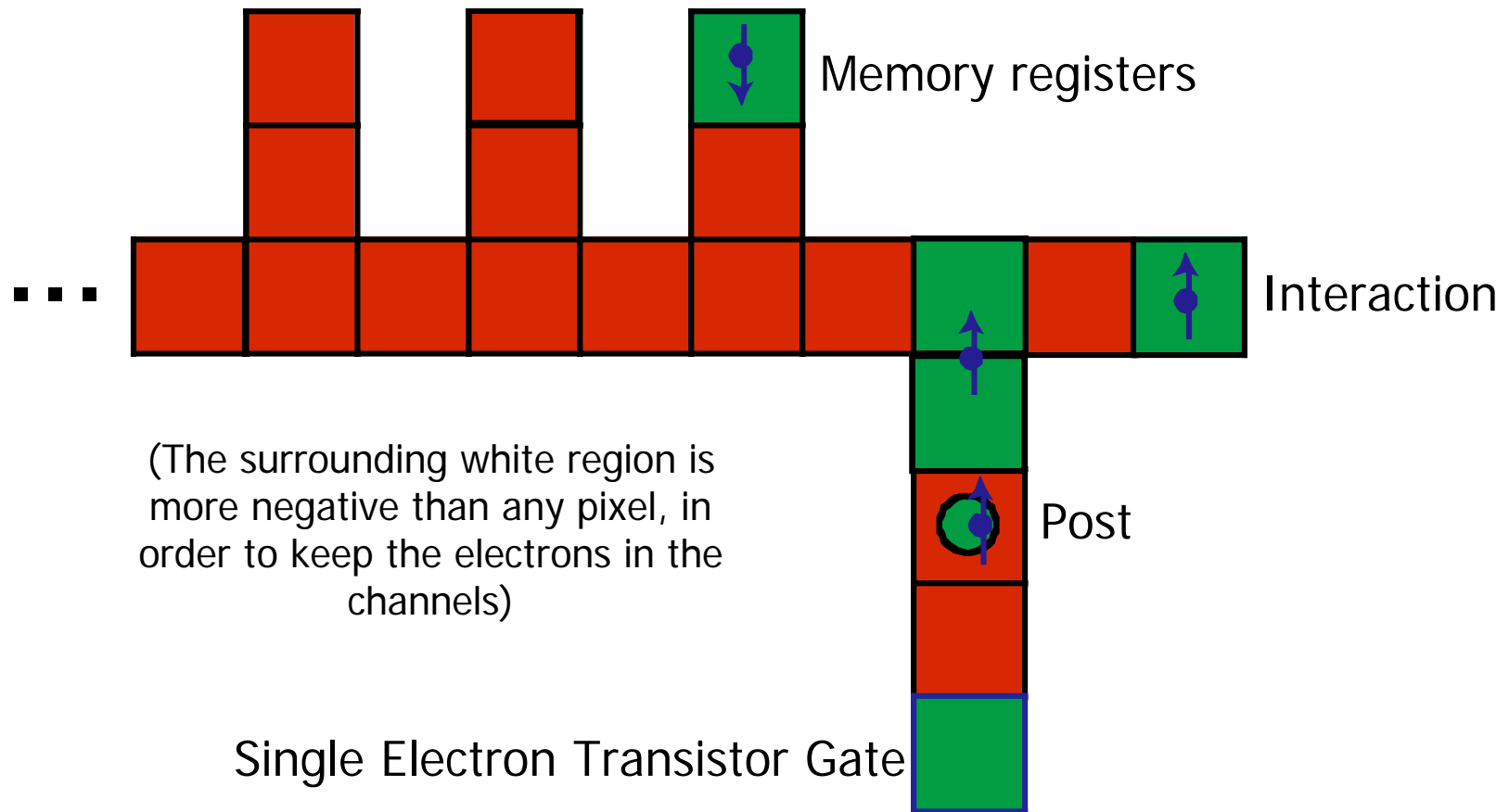
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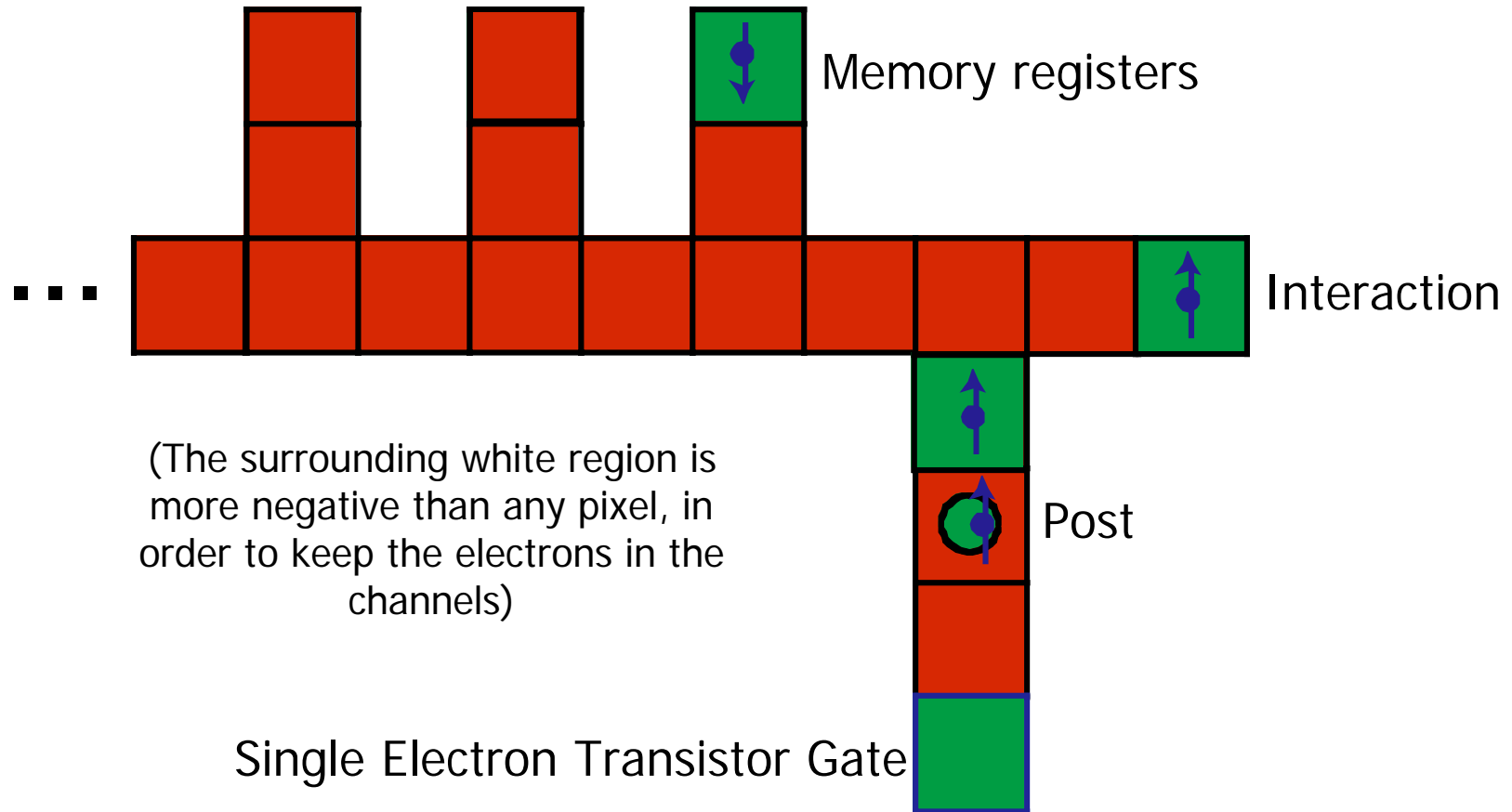
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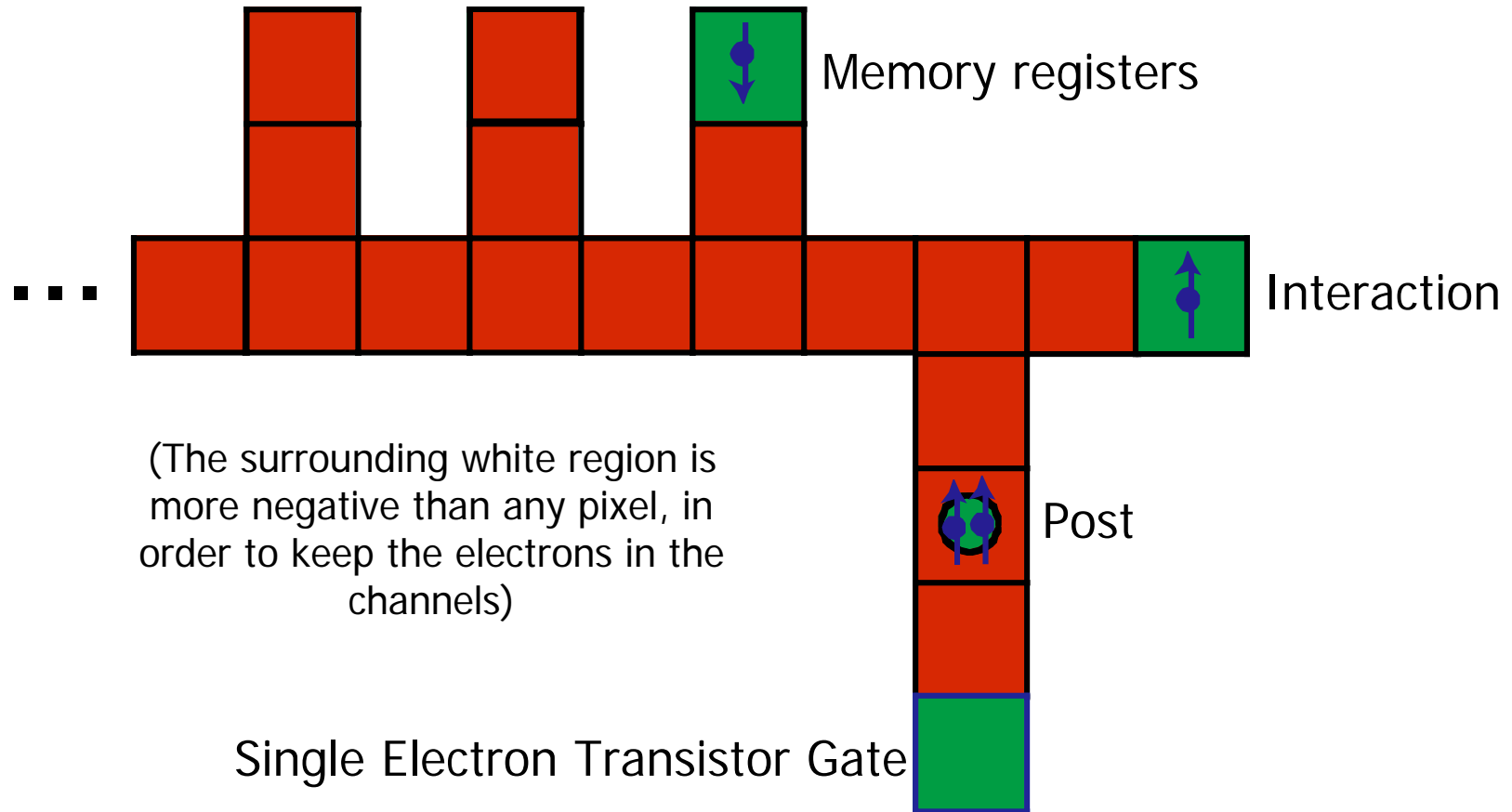
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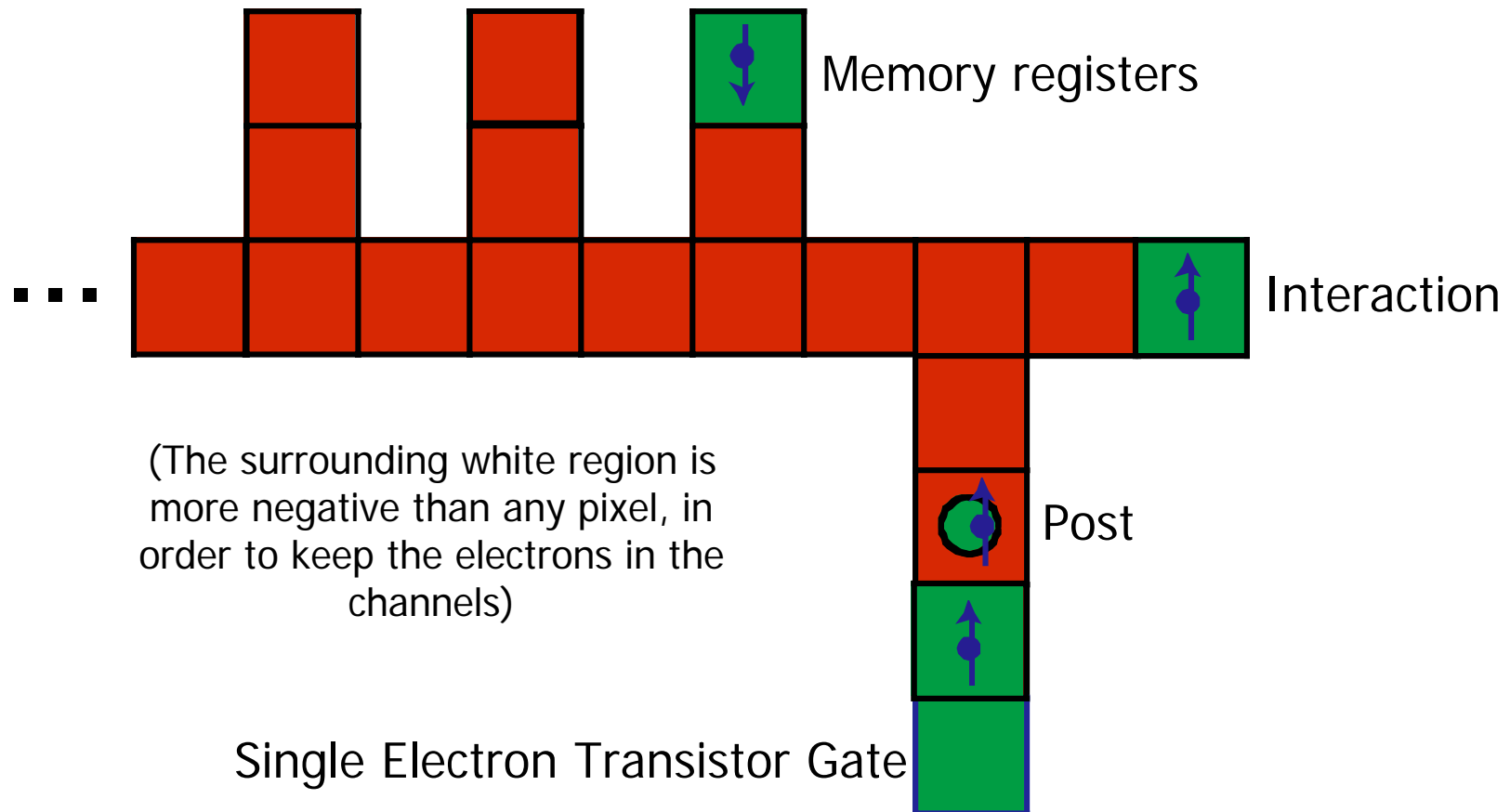
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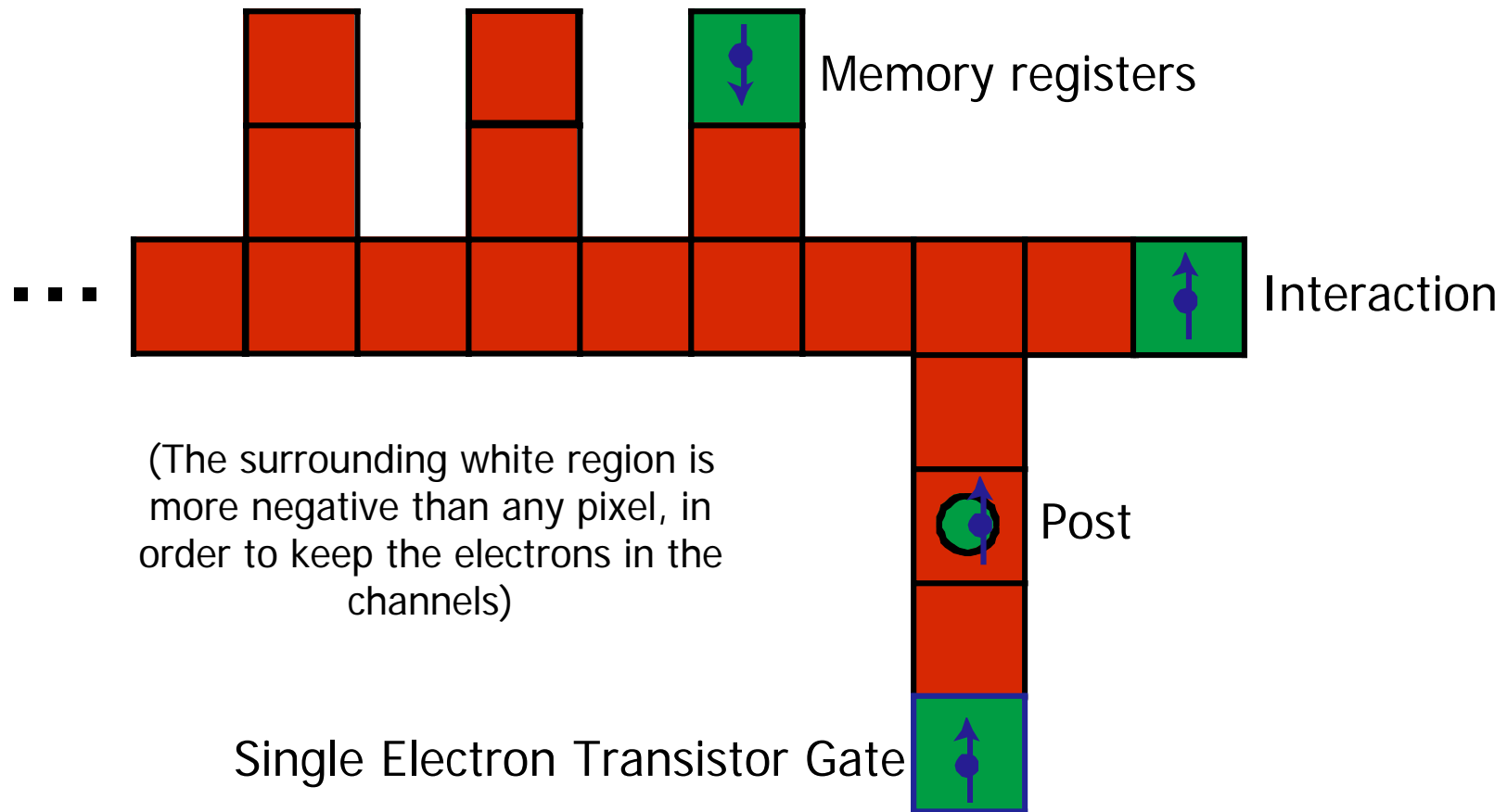
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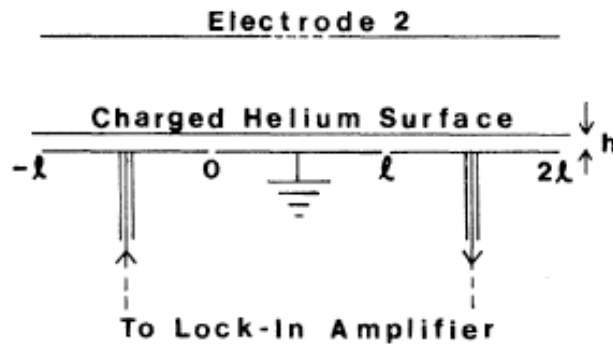
Illustration of Pixel Sequence



- An electron is taken out of memory, brought to interact with another electron and is then read.

Clocking Electrons on Helium

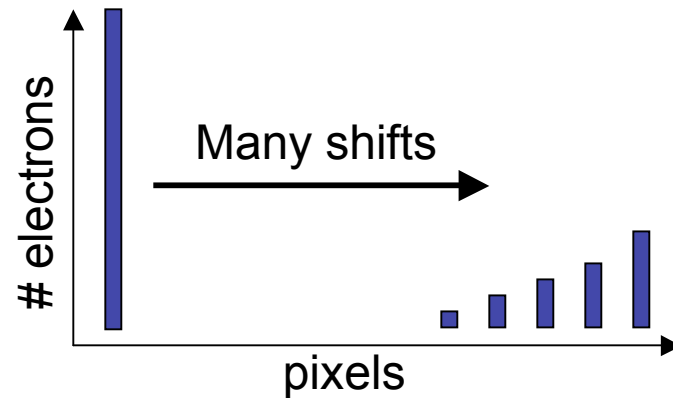
- How readily can we shift electrons around?



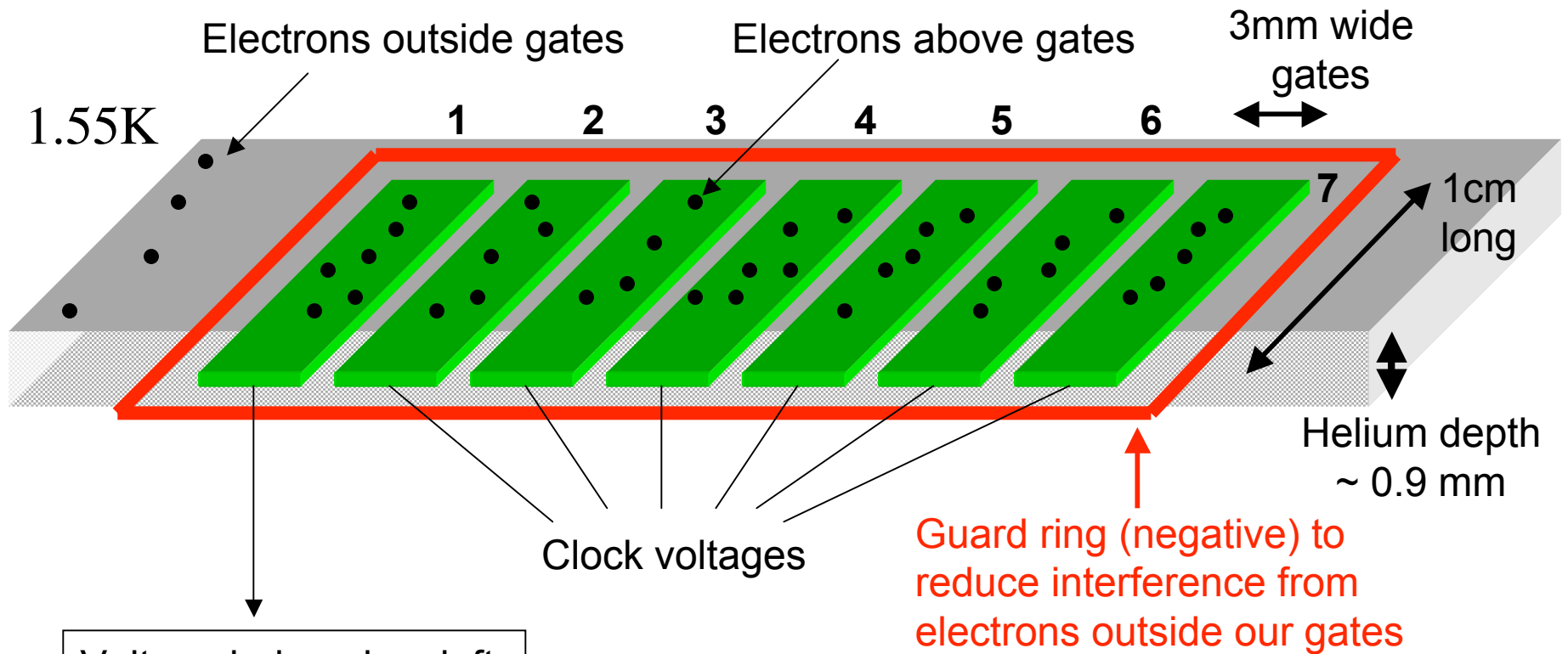
W. T. Sommer and D. J. Tanner,
Physical Review Letters **27**, 20 (1971).

- Measurement of the mobility of the electrons that are free to move.
- These experiments do not tell us the efficiency with which we transfer electrons (CTE = charge transfer efficiency) – charge trapping \Rightarrow CTE \neq 1.0

CTE is normally seen as a smearing of a signal – start with many electrons on one gate and some get left behind with each shift



Thick (~ 0.9 mm) helium device



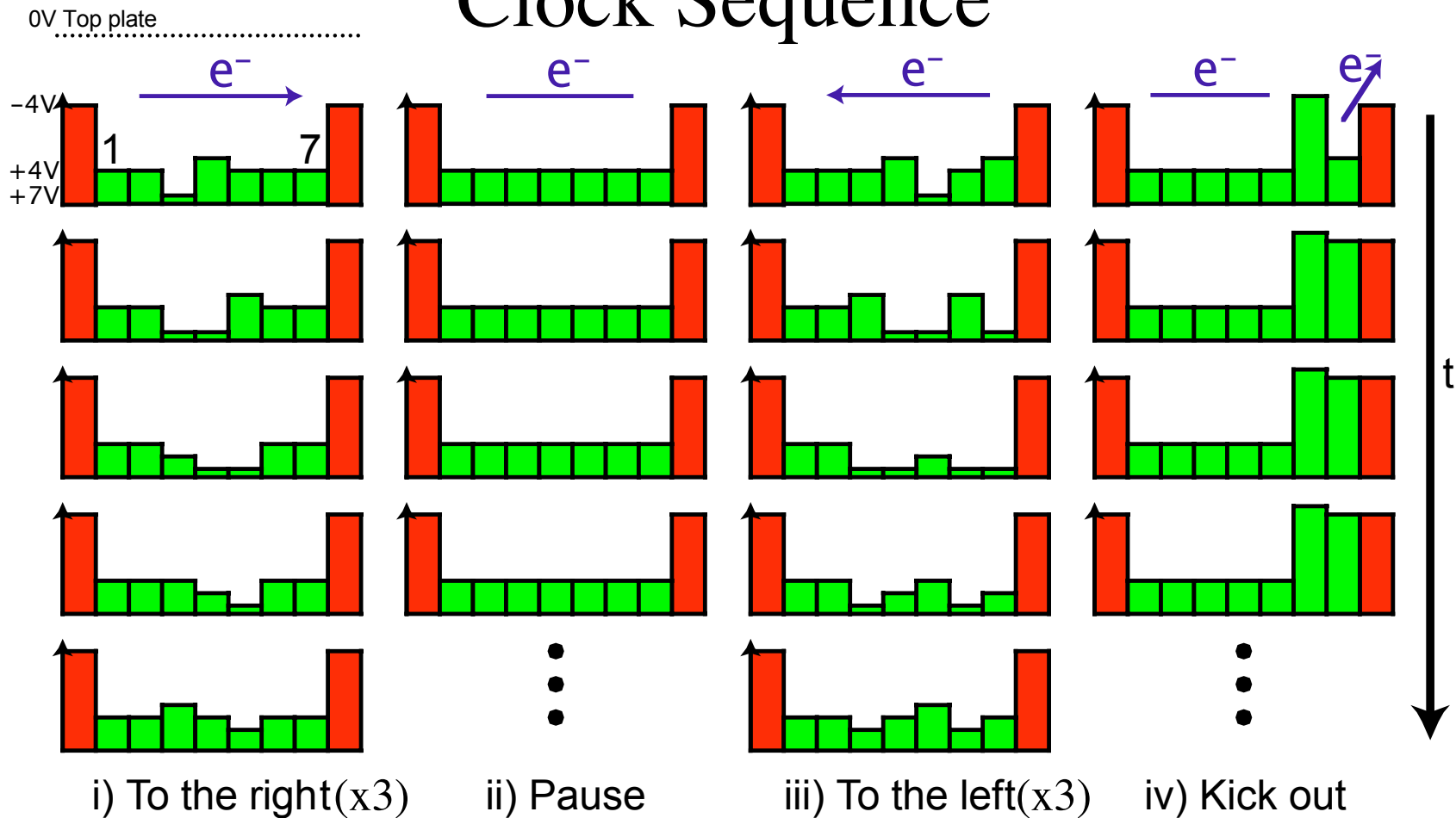
Voltage induced on left gate by the electrons on the helium goes to lockin amplifier

- All gates carry a common dc bias to hold electrons ($+4V \Rightarrow \sim 2.6 \times 10^7$ electrons/cm² or 1 electrons per 4 μm^2)



Picture of the actual device

Clock Sequence

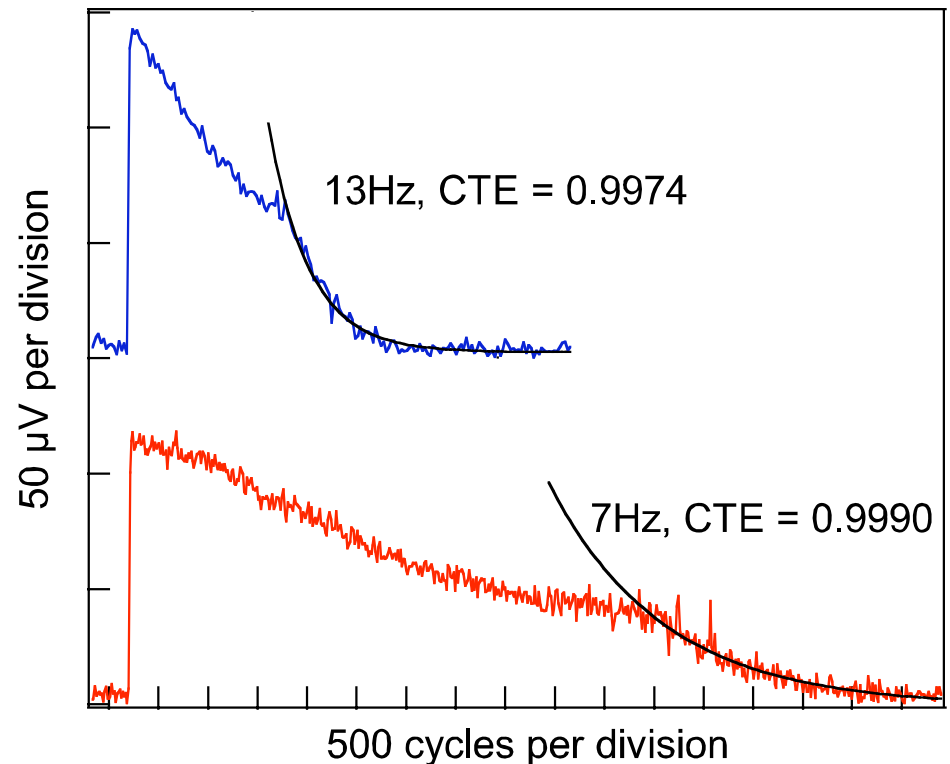


- We don't have long shift register, so use sequence which kicks electrons off the He if they get "stuck" on a gate
 - Clock electrons left and right but after transferring to gate 1 (left), raise energy on gates 6 & 7 (gates negative) to expel remaining electrons from the device
- Variable total period



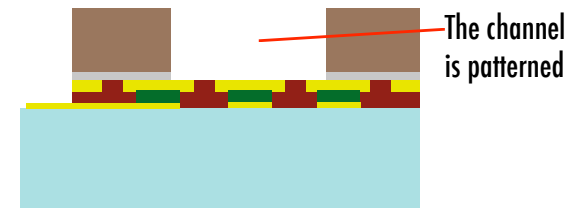
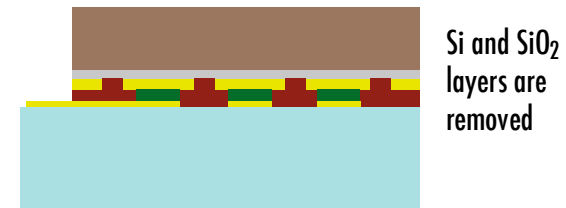
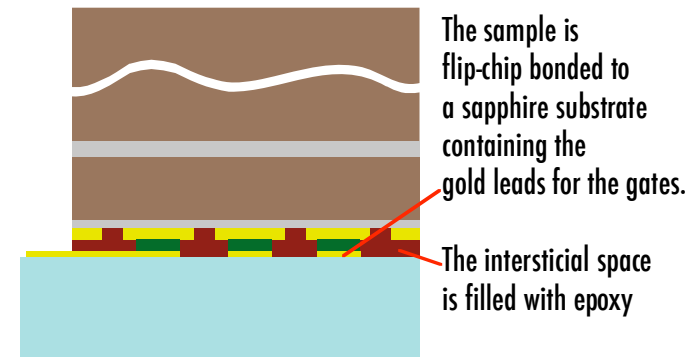
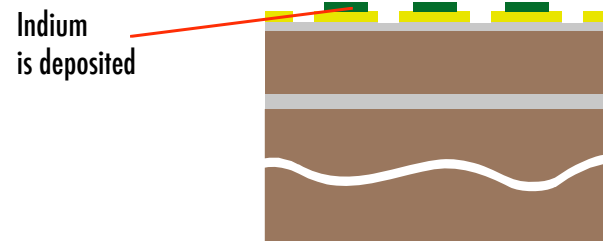
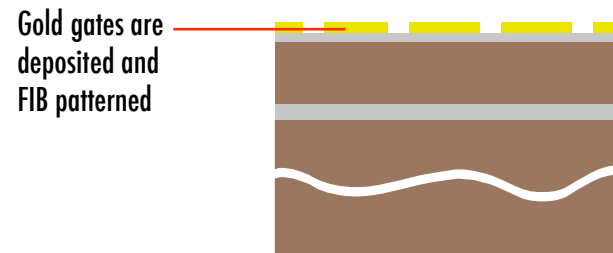
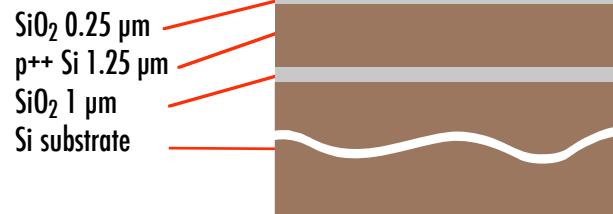
CTE Measurements

- At low frequency (7 Hz) find CTE = .999
- CTE decreases with increasing frequency because electrons must diffuse from one pixel to another, and pixels are huge (3mm)
- CTE would be negligible in Si for our pixel size and electron density

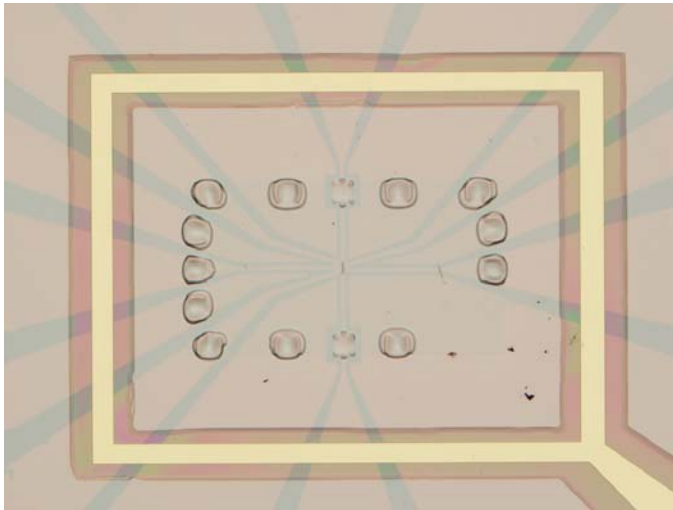


No evidence for strong electron trapping at low frequencies

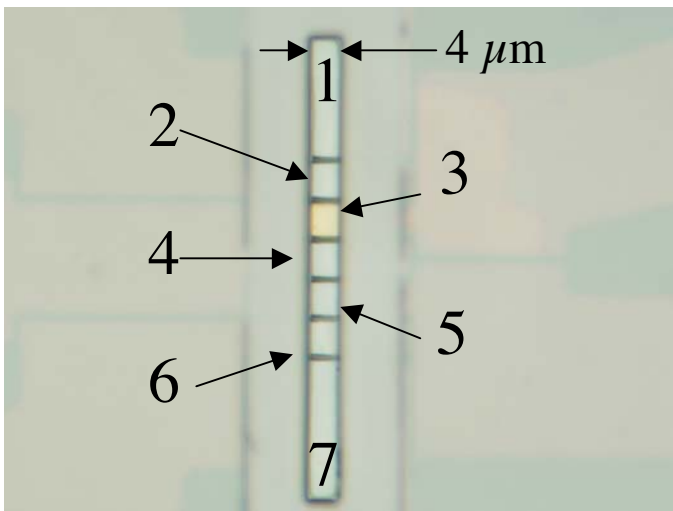
Channel Device Fabrication



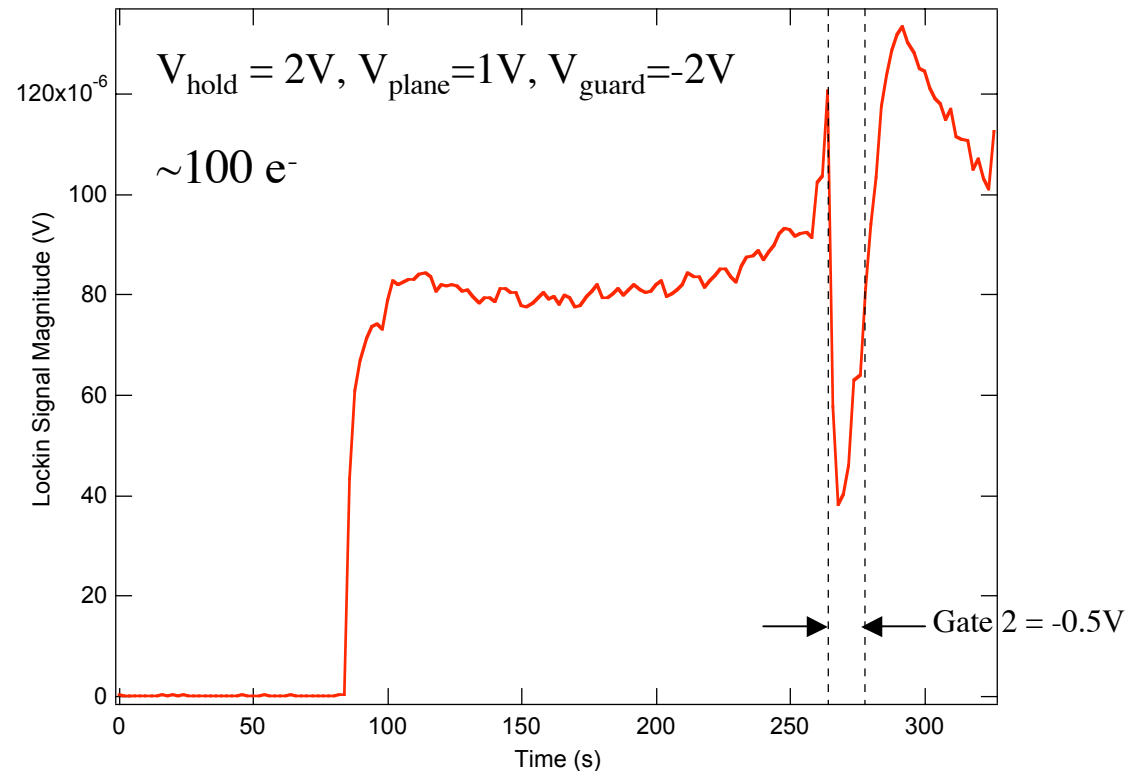
Channel Device ($\sim 1.5\mu\text{m}$)



View of the sample with the guard ring. (The bumps arise from the strain of the flip-chipping)



Zoom on the channel. The $1.5\mu\text{m}$ Silicon layer is open to reveal the underlying metallic gates.



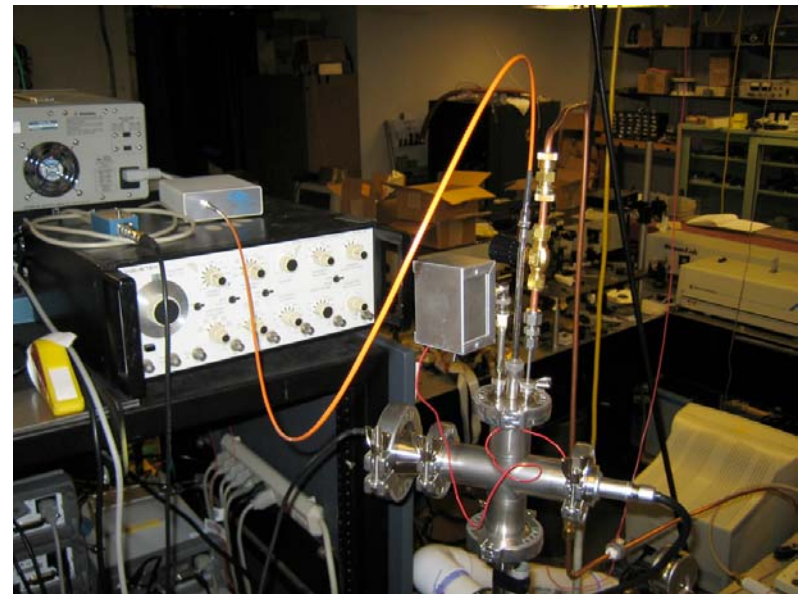
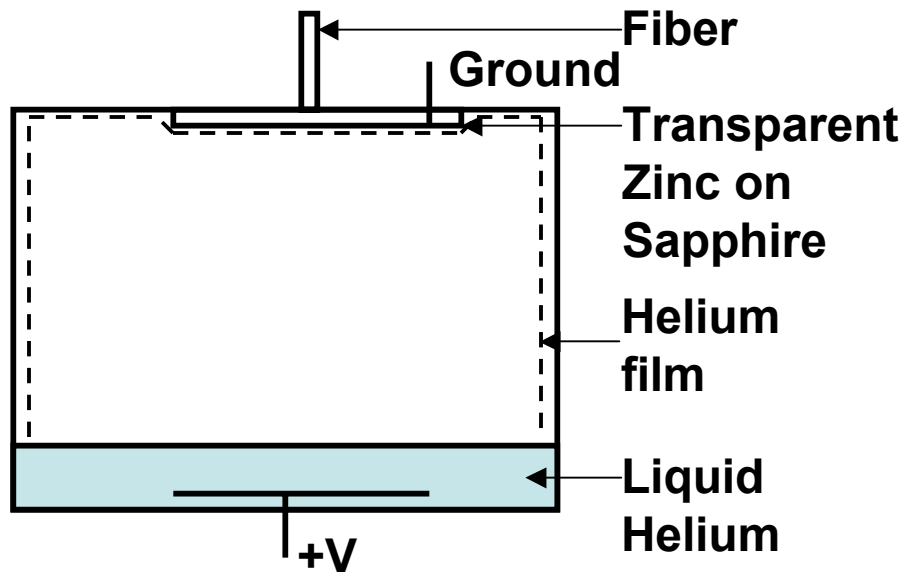
Lockin signal magnitude as a function of time. Making gate number 2 more negative interrupts the signal.

Commercial HEMT at low temperature for readout

Photoemission Source

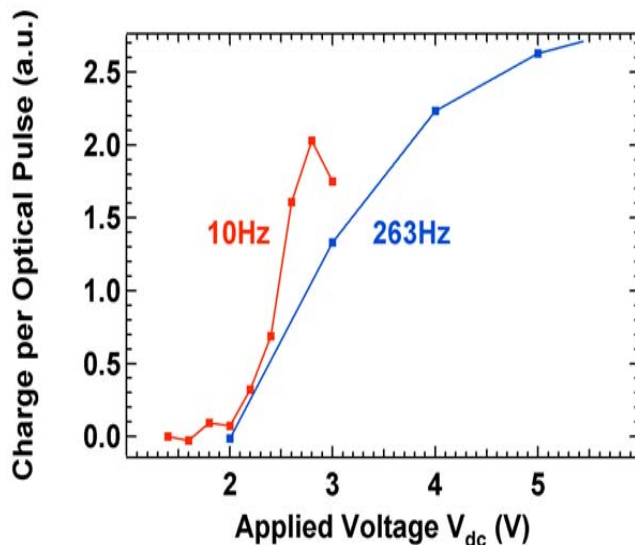
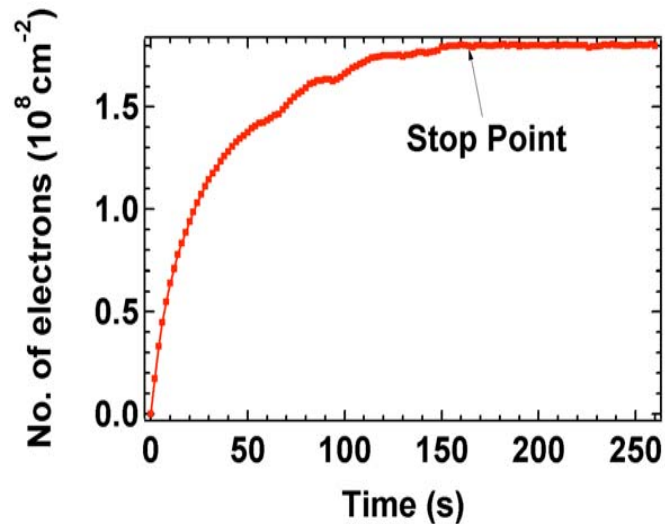
(Shyam Shankar)

- Problem with filaments: too much heat



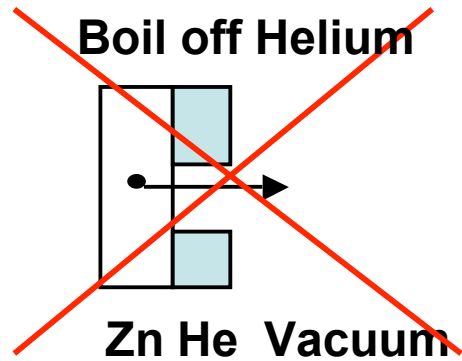
- Wilen and Gianetta (1985)
 - Photoemission from Zinc
 - 1KW arc lamp focused into fiber
- Try low power lamp and small setup

Photoemission Source

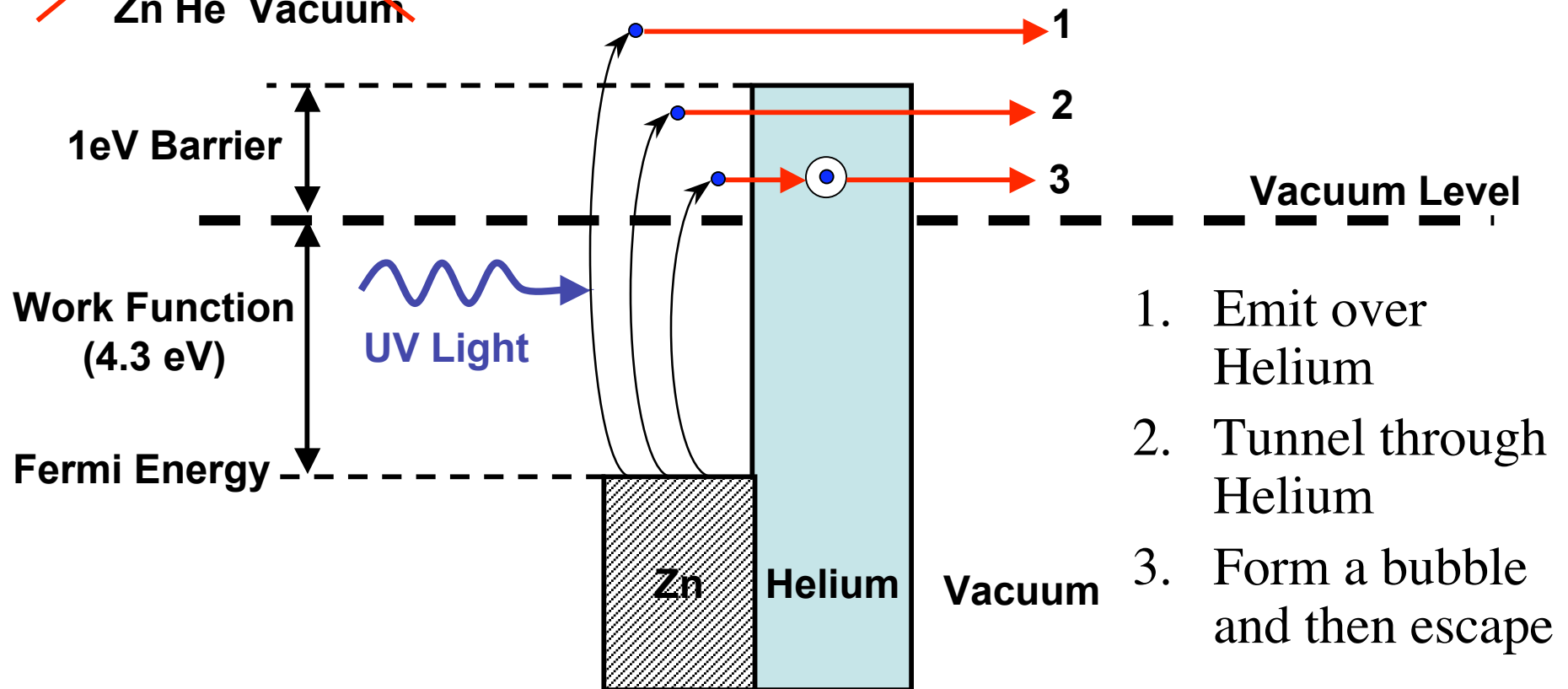


- Pulsed Xenon Source (Ocean Optics PX-2)
- 600 micron Solarization resistant fiber
- 10nm Ti / 20nm Zn on Sapphire substrate
- Get 10^7 - 10^8 electrons in about 20 seconds
- Adjust charge up rate by varying pulse rate
- Find threshold voltage of 2V (4V/cm)

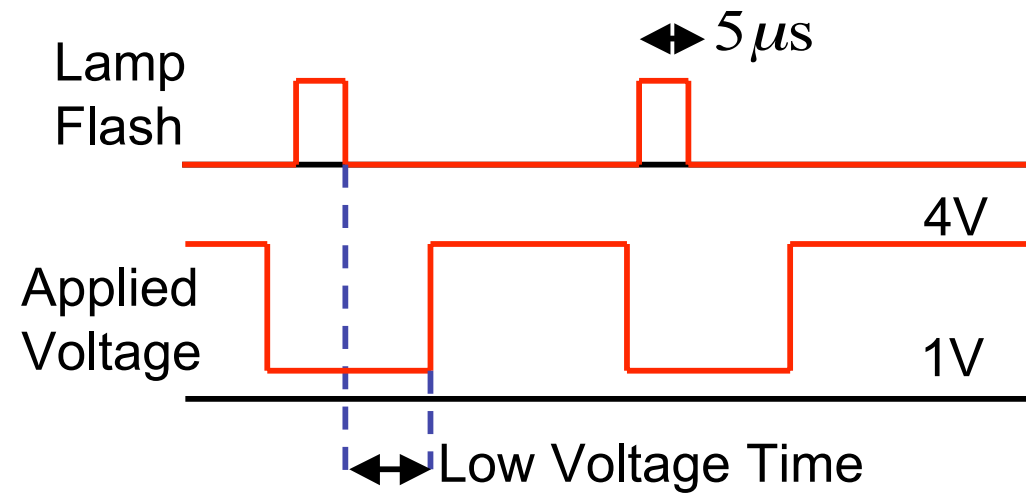
How do the Electrons get through the Helium film?



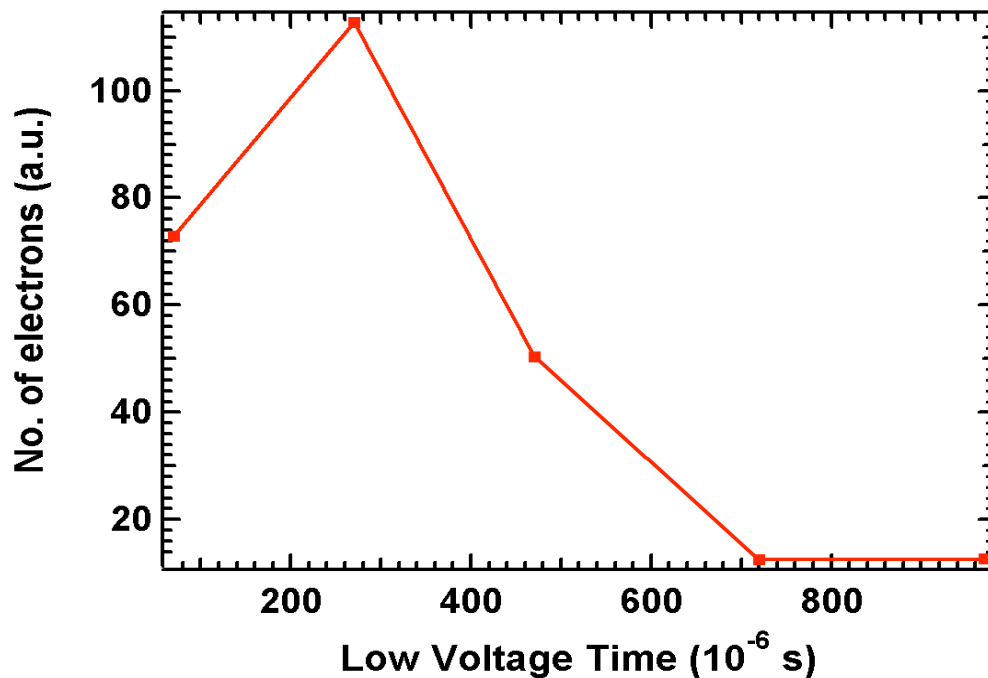
- No – Get 200 times more current if no Helium present
- Only ~100nJ total energy in UV pulse – can't boil film



Bubble States?



- Lamp flashes for $5\mu\text{s}$ every 5ms (200 Hz)
- Low voltage (below 2V threshold) during lamp flash and for a variable time after flash



Some electrons even after almost 1ms

Summary

- Developed a method to measure charge-transfer-efficiency on thick helium without building a long array
 - Measured a of 0.999 at $\sim 10^7$ electrons/cm² and low frequency
 - CTE can be understood in terms of electron diffusion
 - No evidence of electron trapping
 - Cond-mat/0602228
- Demonstrated electron clocking in a 4 μm -wide channel
- Developed an electron photoemission source
 - Small, low power UV source
 - Photoemission mechanism not fully understood